

# The fate of an impurity in a Fermi gas: polaron or molecule?

Pietro Massignan (UAB&ICFO-Barcelona)

in collaboration with: Georg Bruun (Aarhus U.)





# Many-body systems

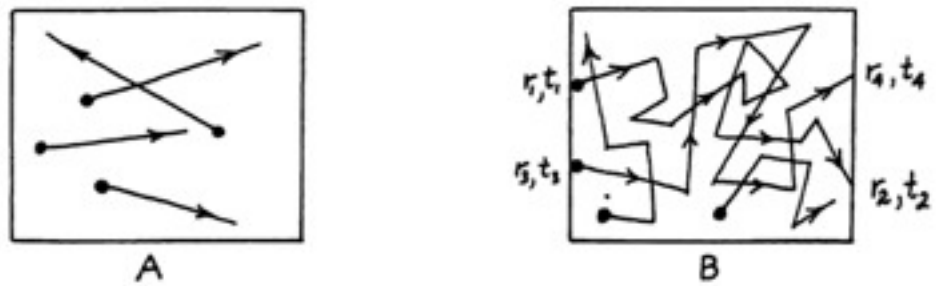


Fig. 0.2 A. Non-interacting Particles  
B. Interacting Particles

(from Richard Mattuck's book)

2

A GUIDE TO FEYNMAN DIAGRAMS [0.0]

Angels on pinhead

Nucleons in nucleus

Electrons in atom

Atoms in molecule

Atoms in solid

Molecules in liquid

Electrons in metal

Fig. 0.1 Some Many-body Systems



# Many-body systems

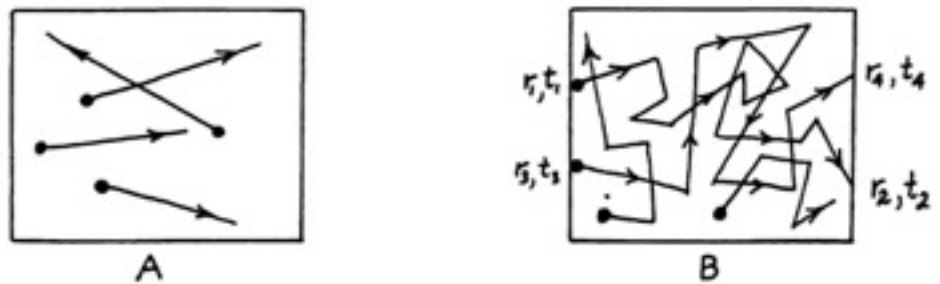


Fig. 0.2 A. Non-interacting Particles  
B. Interacting Particles

2

## A GUIDE TO FEYNMAN DIAGRAMS

[0.0

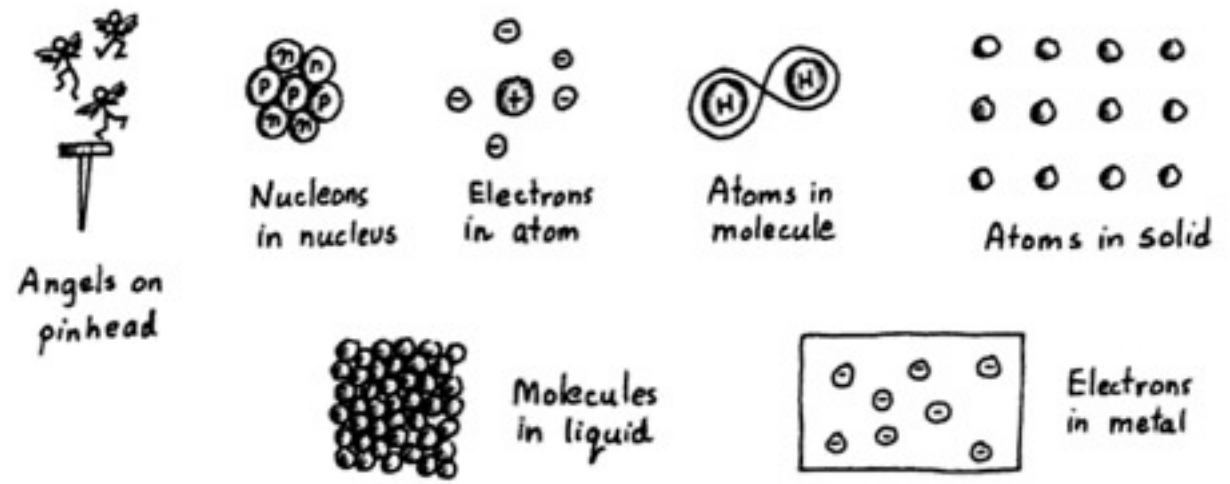


Fig. 0.1 Some Many-body Systems

(from Richard Mattuck's book)

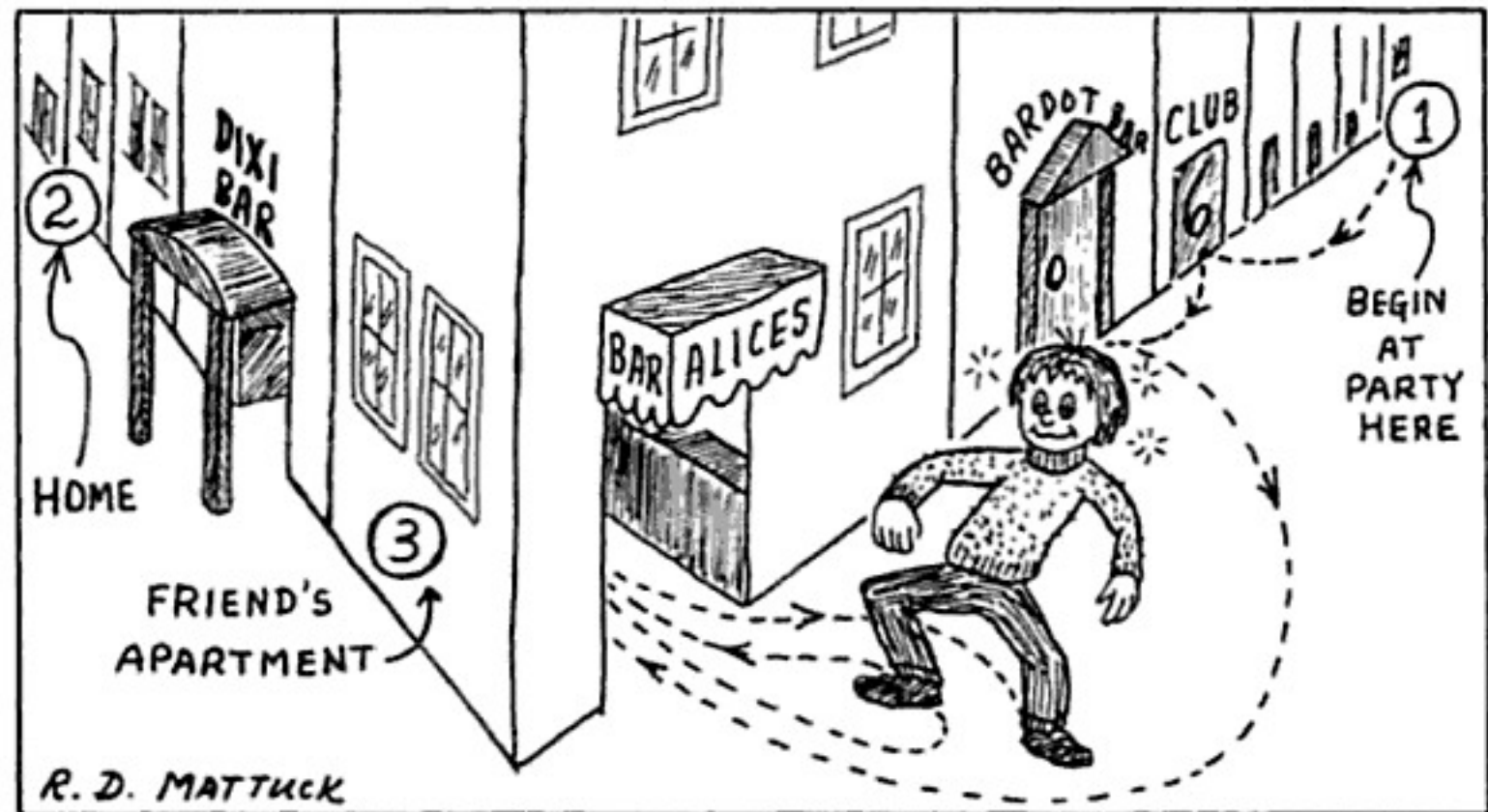


Fig. 1.1 Propagation of Drunken Man



# Quasi-Particles

Landau's idea:  
who cares about real particles?

Excitations behave  
as quasi-particles!

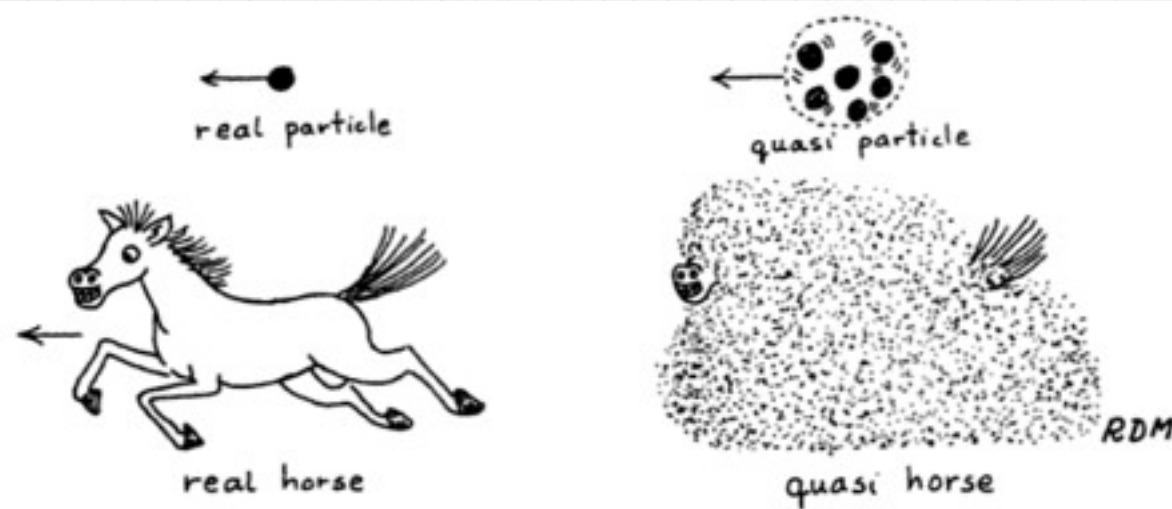
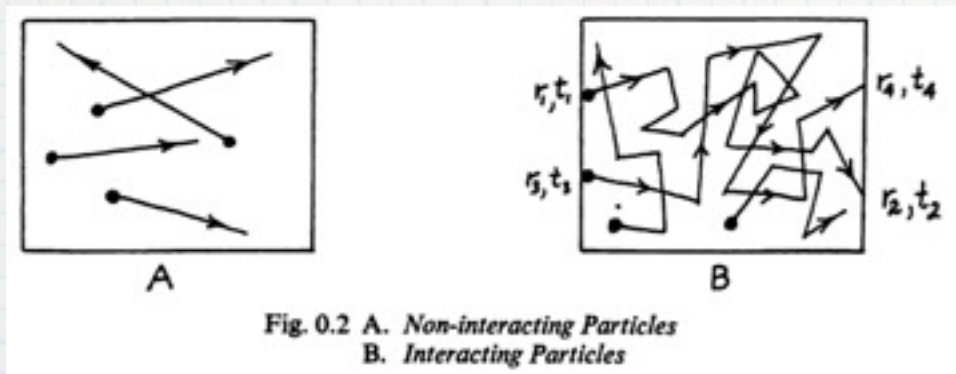


Fig. 0.4 Quasi Particle Concept

- a **QP** is a "free particle" with:
- @ renormalized mass
  - @ self-energy
  - @ shielded interactions
  - @ lifetime

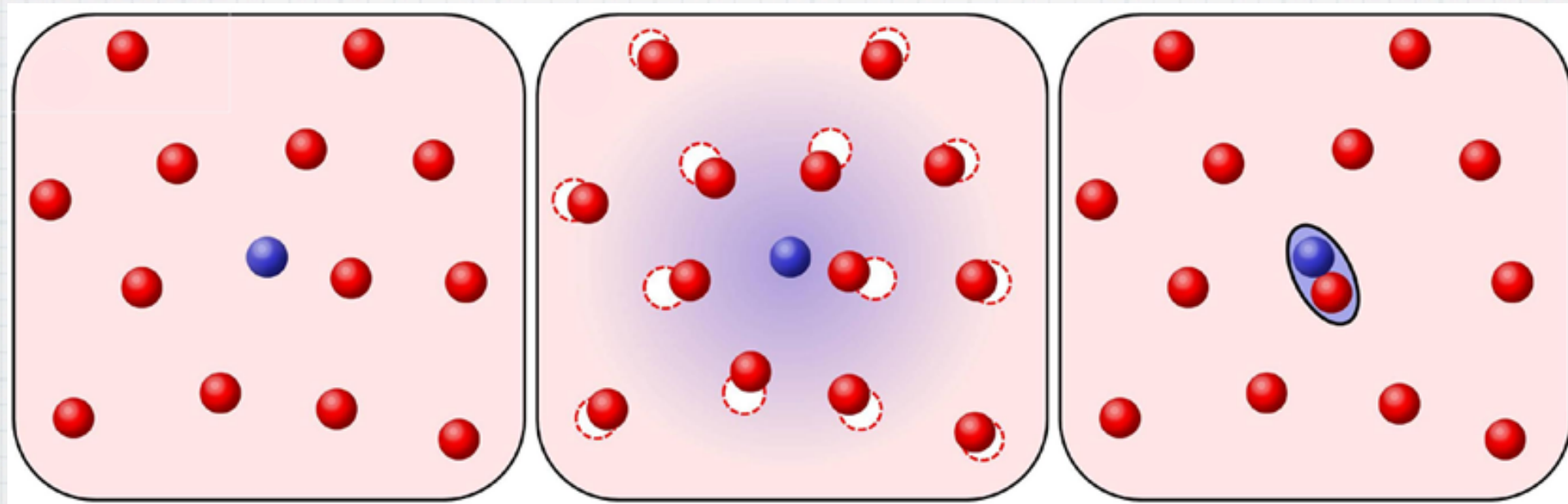


# The impurity problem

- non-interacting ↑ Fermi sea ( $N \gg 1$ )
- a single ↓ impurity

BCS  $\xrightarrow{\text{Attraction strength}}$  BEC

$(k_{Fa})^{-1} < 0$



$(k_{Fa})^{-1} > 0$

free particle

QP (polaron)

molecule



# Polaron Ansatz

(F. Chevy, PRA 2006)

the ↓ impurity

$$|\psi_{\mathbf{p}}\rangle = \phi_0 |\mathbf{p}\rangle_{\downarrow} |0\rangle_{\uparrow} + \sum_{q < k_F} \phi_{\mathbf{q}\mathbf{k}} |\mathbf{p} + \mathbf{q} - \mathbf{k}\rangle_{\downarrow} c_{\mathbf{k}\uparrow}^{\dagger} c_{\mathbf{q}\uparrow} |0\rangle_{\uparrow}$$

non-interacting ↑ Fermi sea

Particle-Hole dressing

This Ansatz gives a very good agreement with MC results for the energy and  $m^*$ , even at unitarity.

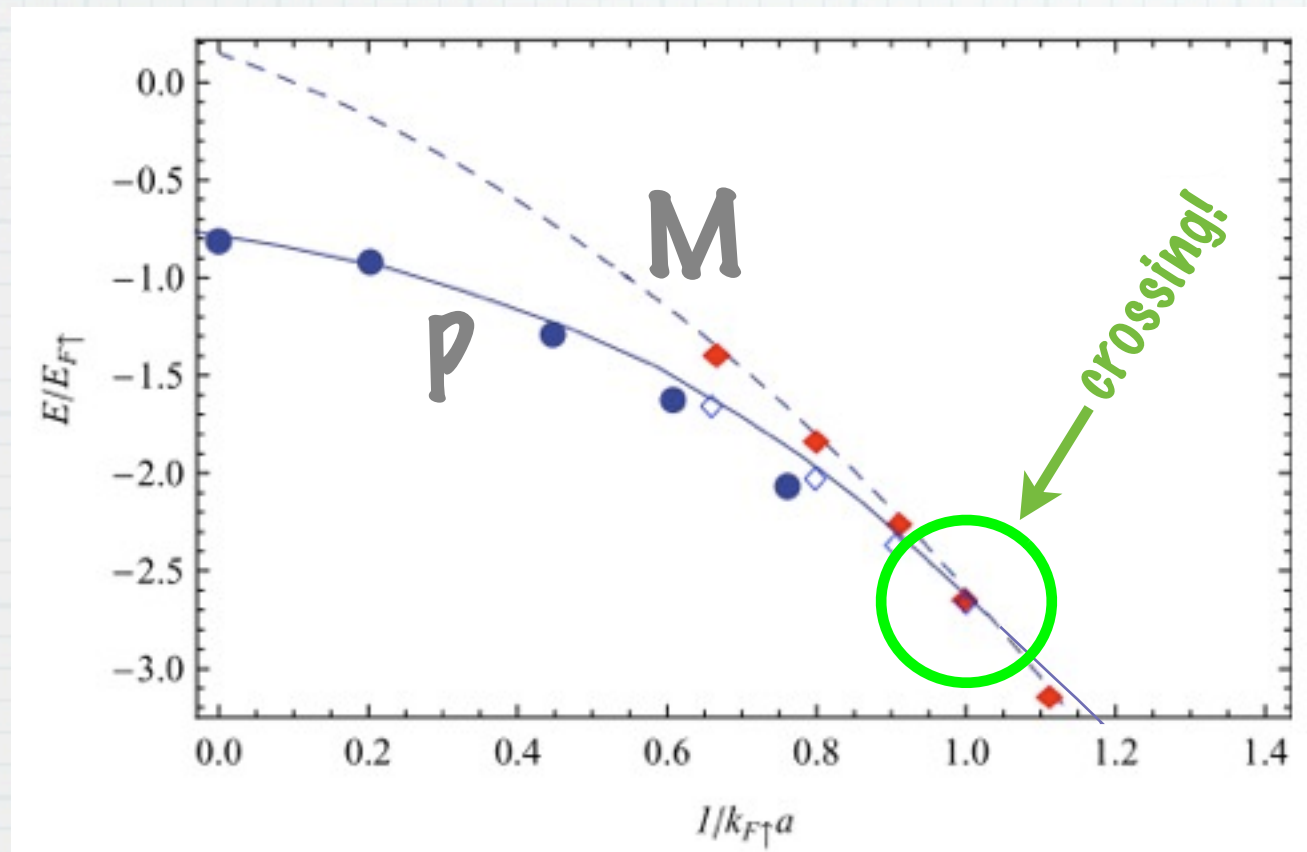
The variational treatment has a diagrammatic equivalent. It corresponds to the forward scattering, or ladder, approximation.

(Combescot et al., PRL 2007)

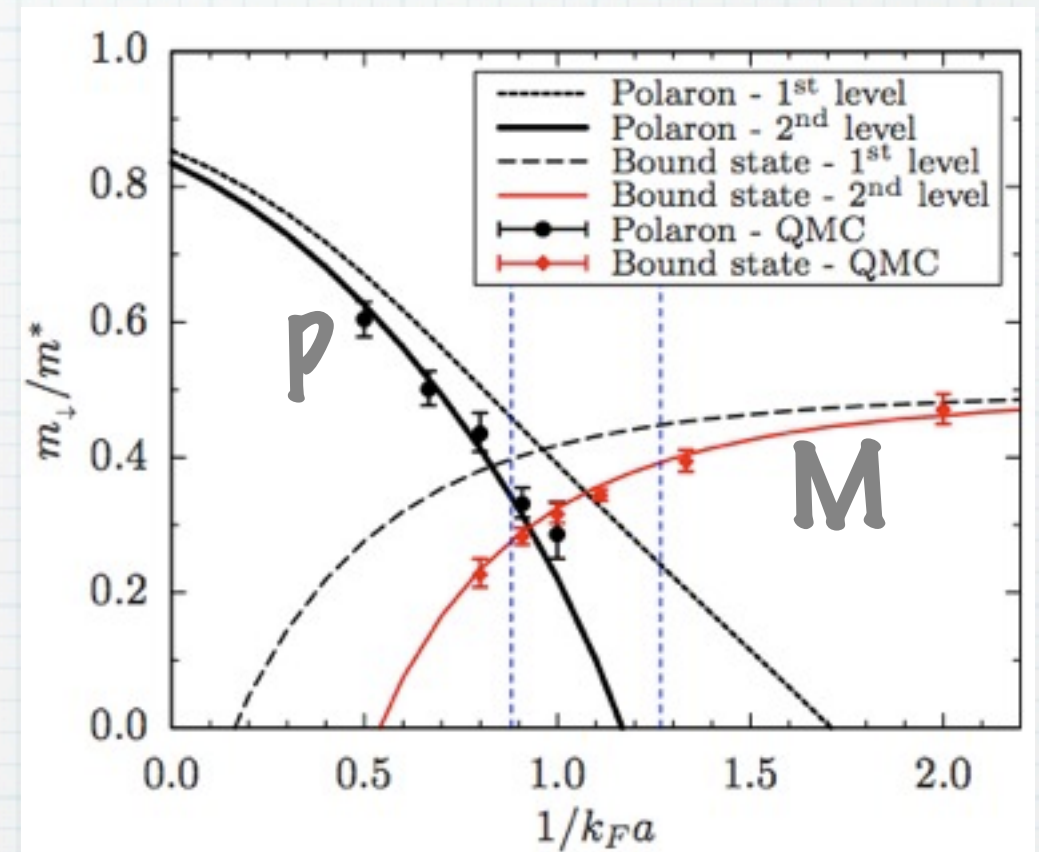


# QP parameters

## Self-energy $\Sigma$



## Effective mass $m^*$



Theory (QMC, variational, diagrammatic): Prokof'ev&Svistunov, Chevy, Recati, Lobo, Stringari, Combescot, Leyronas, Massignan&Bruun, Zwerger, Punk, Stoof, Mora,...

Experiments: MIT, ENS

**P-P Interactions:** Mora&Chevy, PRL 2010; Zhenhua, Zöllner&Pethick, arXiv:1006.4723

**P&M lifetimes:** Bruun&Massignan, PRL 2010



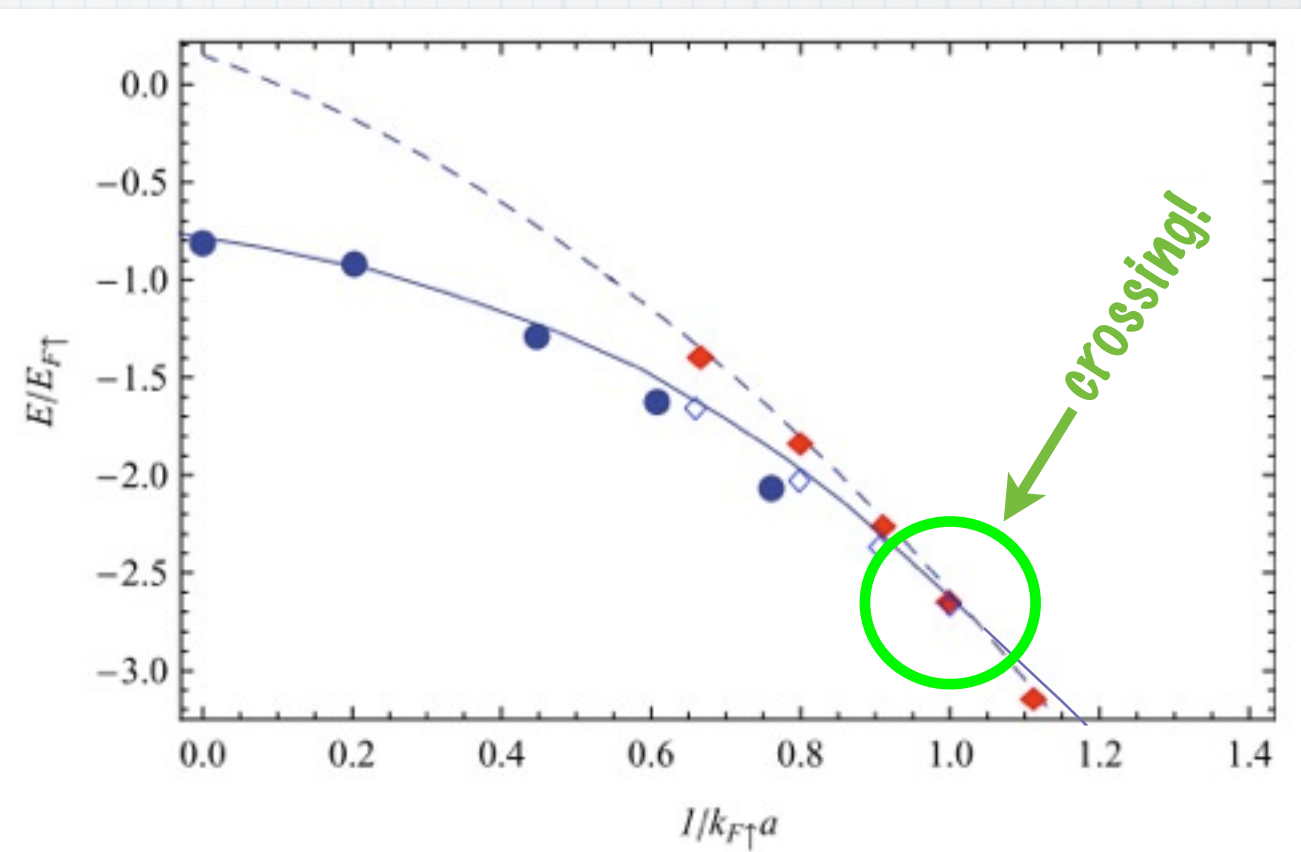
# Bottom line: long life to quasi-particles!

$$\Gamma_P \sim Z_M(k_F a) (m_M^*)^{3/2} (\Delta\omega)^{9/2}$$

$$\Gamma_M \sim Z_P(k_F a) (m_P^*)^{3/2} (|\Delta\omega|)^{9/2}$$

$$\Delta\omega = \omega_P - \omega_M$$

Long lifetimes  
~ 10-100ms



Bruun&Massignan, PRL 2010



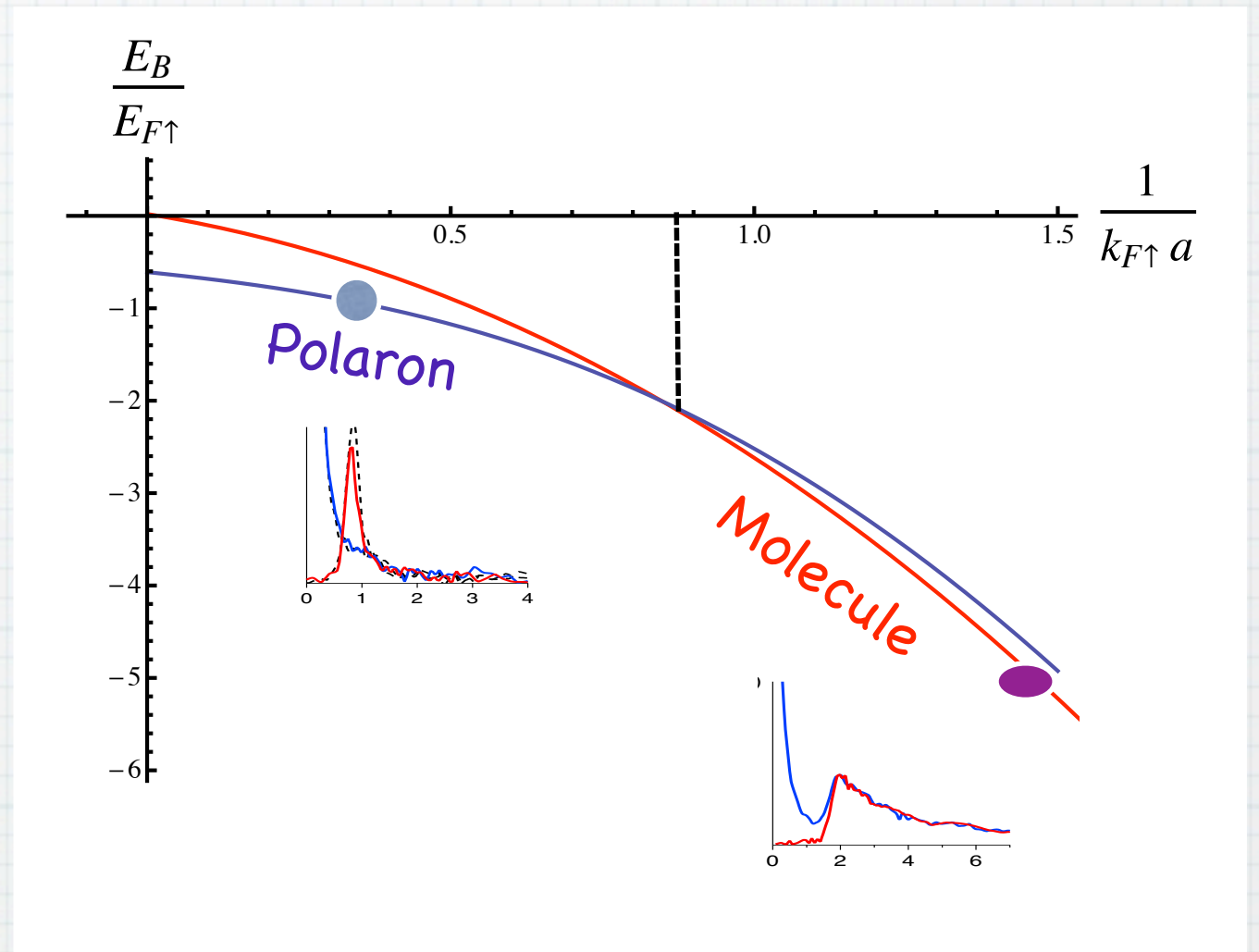
# Experimental observation

## Methods:

- RF spectra
- Collective modes to measure  $m^*$  vs. time

## Issues:

- \* No decay to deeply bound molecular states
- \* Phase separation?
  - \* stabilized by finite  $T$
  - \* work with  $m_{\downarrow} \neq m_{\uparrow}$
  - \* use bosonic impurities





# Pol $\rightarrow$ Mol decay

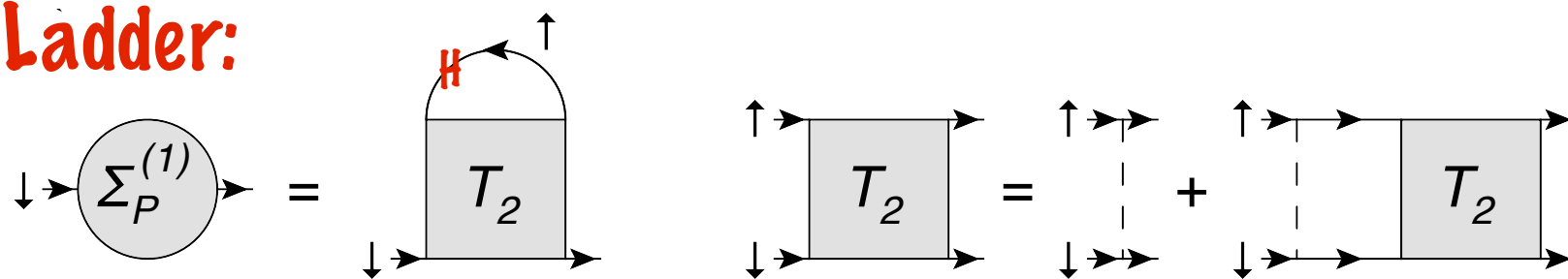
$$\Delta\omega = \omega_P - \omega_M > 0$$

Polaron:  $G_{\downarrow}(\mathbf{p}, z)^{-1} = G_{\downarrow}^0(\mathbf{p}, z)^{-1} - \Sigma_P(\mathbf{p}, z)$

Decay rate:  $\Gamma_P = -\text{Im}\Sigma_P(p=0, \omega_P)$

Hole expansion:  $\Sigma_P(\mathbf{p}, z) = \Sigma_P^{(1)}(\mathbf{p}, z) + \Sigma_P^{(2)}(\mathbf{p}, z) + \dots$

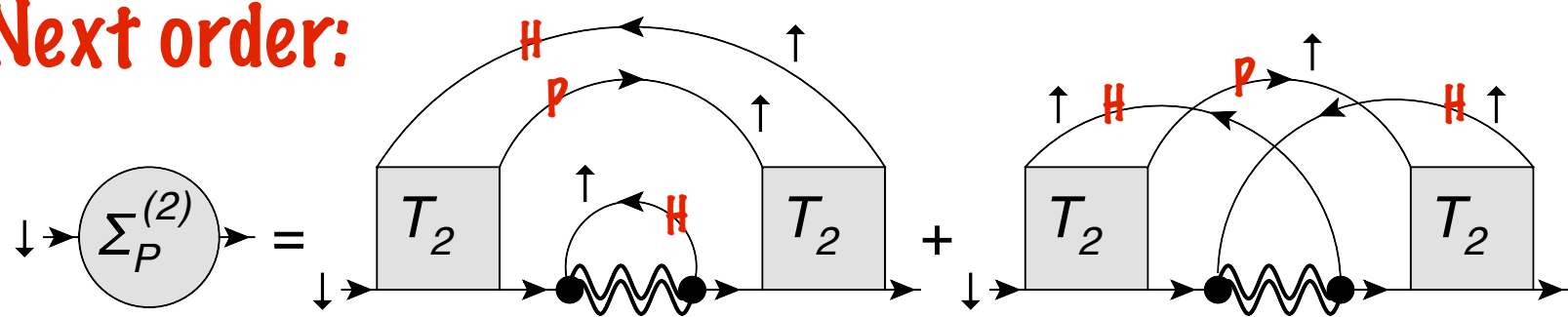
**Ladder:**



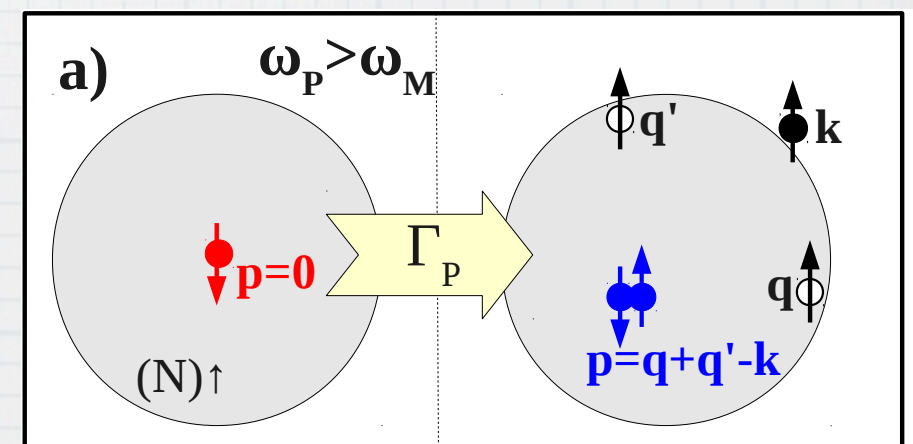
**no damping  
in the ladder approx.**

**3-body process**

**Next order:**



**dressed molecule**





molecule w.f.  
in vacuum:

$$\phi_q = \frac{\sqrt{8\pi a^3}}{1 + q^2 a^2} \quad \text{or} \quad \phi_r \propto \frac{e^{-r/a}}{r}$$

dressed molecule:



$$D(\mathbf{p}, \omega) \simeq \frac{Z_M}{\omega - \omega_M - p^2/2m_M^*}$$

atom-molecule  
coupling:

$$\bullet = \frac{1}{g(\mathbf{p}, z)} = \int \frac{d^3 q}{(2\pi)^3} \phi_q \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} \sim -\sqrt{\frac{m_r^2 a}{2\pi}}$$

(Bruun&Pethick, PRL 2004)

$$\Gamma_P = \frac{g^2 Z_M}{2} \int \frac{d^3 k d^3 q d^3 q'}{(2\pi)^9} [F(\mathbf{q}, \mathbf{k}, \omega_P) - F(\mathbf{q}', \mathbf{k}, \omega_P)]^2 \delta \left( \Delta\omega + \xi_{\mathbf{q}\uparrow} + \xi_{\mathbf{q}'\uparrow} - \xi_{\mathbf{k}\uparrow} - \frac{(\mathbf{q} + \mathbf{q}' - \mathbf{k})^2}{2m_M^*} \right)$$

$$q, q' < k_F, \quad k > k_F$$

$$F(\mathbf{q}, \mathbf{k}, \omega) = T_2(\mathbf{q}, \omega + \xi_{\mathbf{q}\uparrow}) G_{\downarrow}^0(\mathbf{q} - \mathbf{k}, \omega + \xi_{\mathbf{q}\uparrow} - \xi_{\mathbf{k}\uparrow})$$



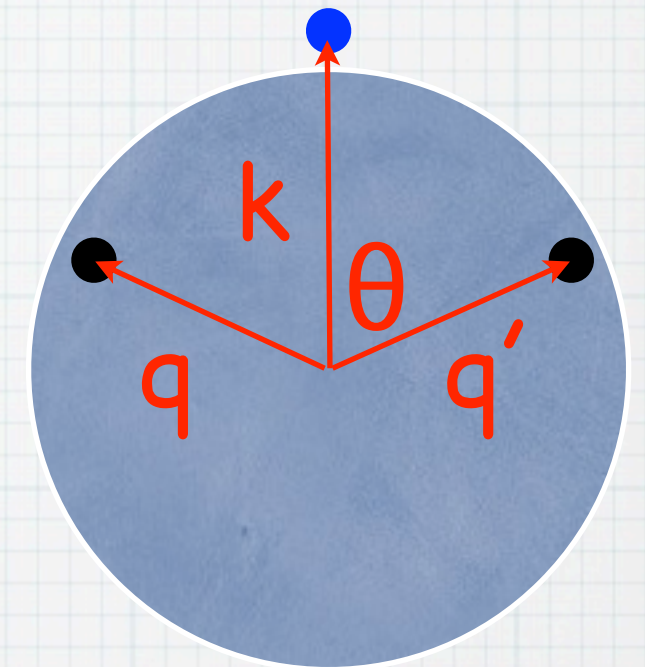
In the neighborhood of the P-M crossing,

$$\Delta\omega \ll \epsilon_F$$

$$q \simeq k \simeq k' \simeq k_F$$

$$\int \frac{d^3k d^3q d^3q'}{(2\pi)^9} \delta(\dots) \sim (m_M^*)^{3/2} (\Delta\omega)^{7/2}$$

The P+H+H form an equilateral triangle, since  $q + q' - k \sim 0$



At the crossing, Fermi antisymmetry yields a vanishing of the matrix element!

$$F(\mathbf{q}, \mathbf{k}, \omega_P) - F(\mathbf{q}', \mathbf{k}, \omega_P)$$

the angular dependence of F is only on  $\theta$

Expand matrix element around the equilateral shape to get an extra factor of  $\Delta\omega$ :

$$\Gamma_P \sim Z_M(k_F a) (m_M^*)^{3/2} (\Delta\omega)^{9/2}$$

**1<sup>st</sup> order transition between the P&M states (no coupling at the crossing)**

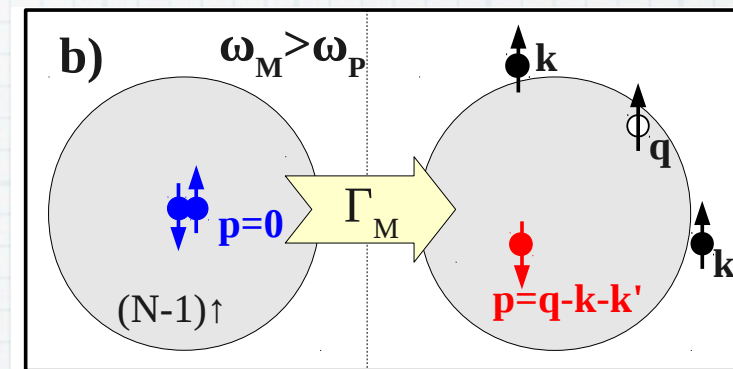


# Mol $\rightarrow$ Pol decay

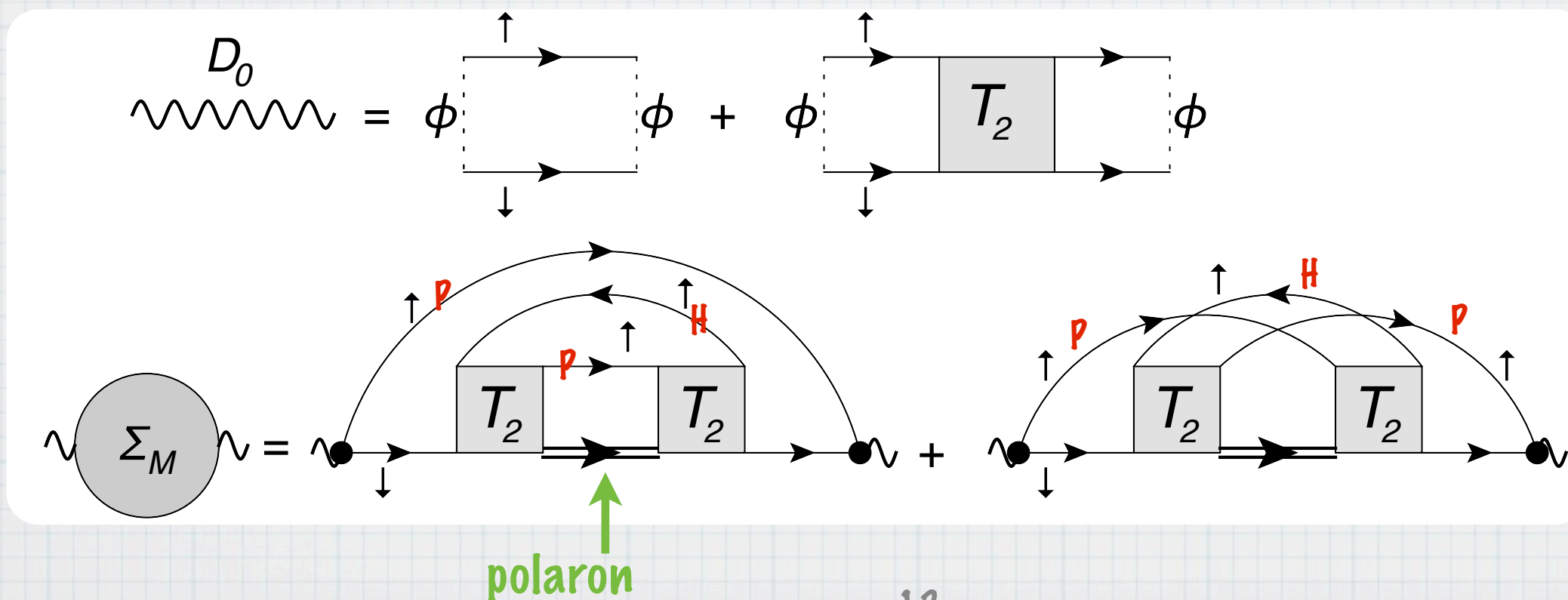
$$\Delta\omega = \omega_P - \omega_M < 0$$

**Molecule:**  $D(\mathbf{p}, z)^{-1} = D_0(\mathbf{p}, z)^{-1} - \Sigma_M(\mathbf{p}, z)$

**Decay rate:**  $\Gamma_M = -\text{Im}\Sigma_M(p=0, \omega_M)$



**Vacuum:**  $D_0(\mathbf{p}, z) = \int d^3\check{q} \phi_q^2 \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} + \frac{T_2(\mathbf{p}, z)}{g(\mathbf{p}, z)^2}$



3-body process



$$\Gamma_M = \frac{g^2 Z_P}{2} \int \frac{d^3 k d^3 k' d^3 q}{(2\pi)^9} [C(\mathbf{q}, \mathbf{k}, \omega_M) - C(\mathbf{q}, \mathbf{k}', \omega_M)]^2 \delta \left( |\Delta\omega| + \xi_{q\uparrow} - \xi_{k\uparrow} - \xi_{k'\uparrow} - \frac{(\mathbf{q} - \mathbf{k} - \mathbf{k}')^2}{2m_P^*} \right)$$

In the neighborhood of the M-P crossing,  $\Gamma_M \sim Z_P(k_F a) (m_P^*)^{3/2} (|\Delta\omega|)^{9/2}$

For both decay processes,  
very **long lifetimes** are ensured by:

- limited phase-space
- Fermi antisymmetry

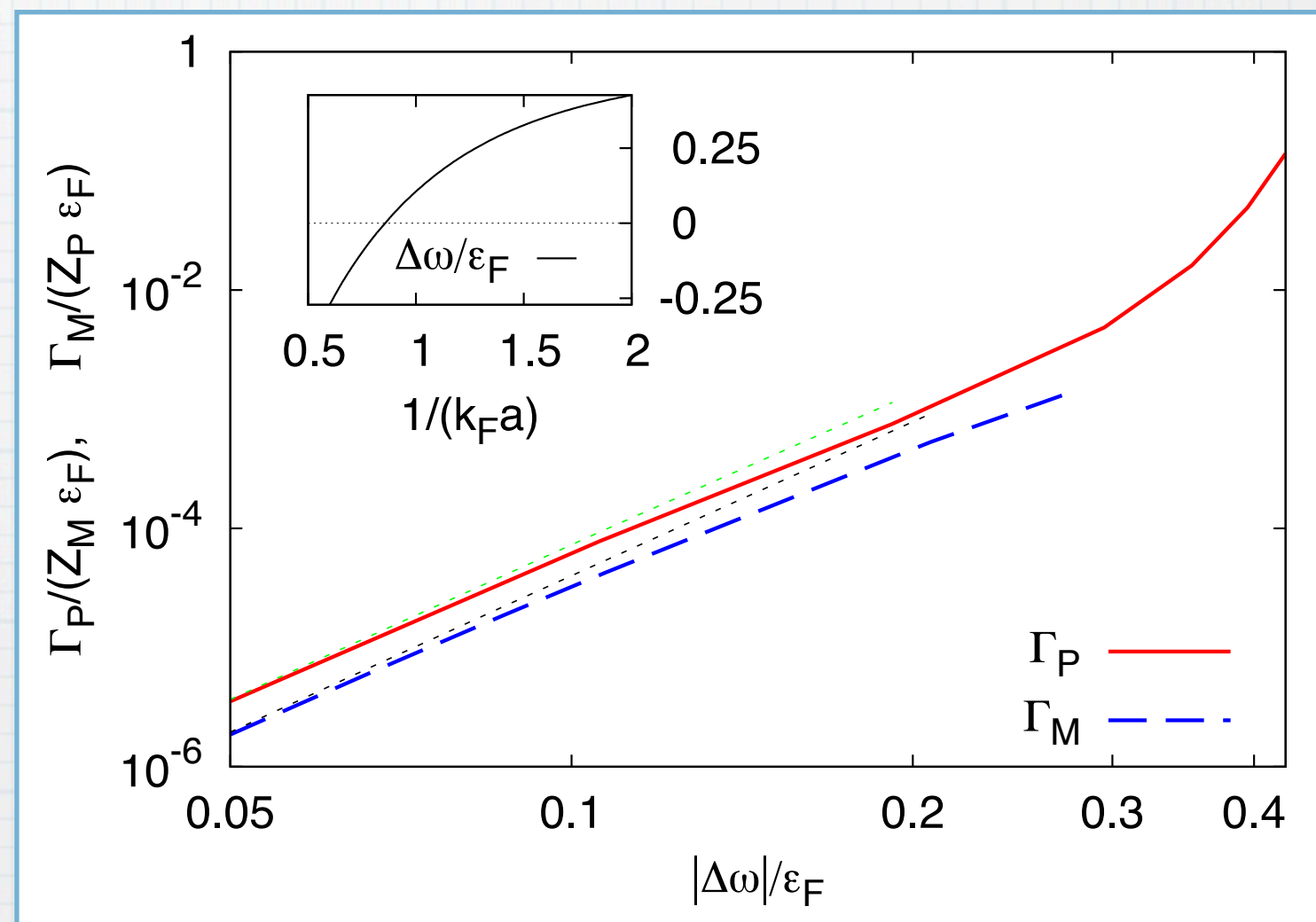
**much longer than usual Fermi liquids**

In the numerics:

$$\omega_M = -\frac{\hbar^2}{2m_r a^2} - \epsilon_F + g_3 n_\uparrow$$

$$a_3 = 1.18a$$

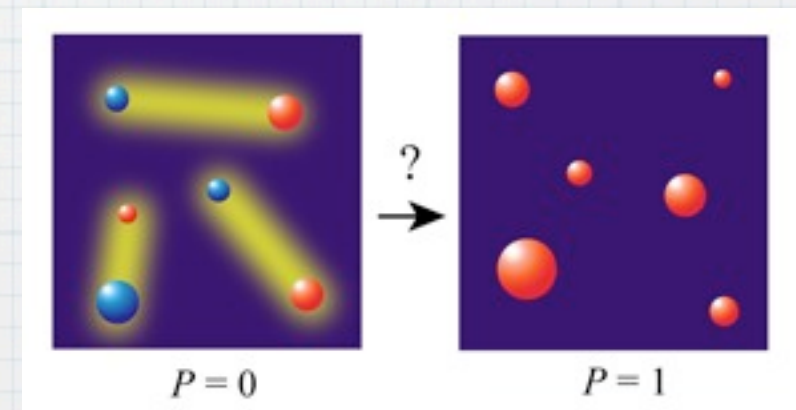
$$T_2(\mathbf{p}, \omega) = \frac{2\pi a/m_r}{1 - \sqrt{2m_r a^2 \left( \frac{p^2}{2m_M} - \omega - \epsilon_F + g_3 n_\uparrow \right)}}$$





# Conclusions

- The impurity problem contains a sharp Polaron-Molecule transition
- At small momenta, the process coupling molecules and polarons requires at least 3-bodies
- The P-M decay is strongly suppressed due to a combination of small final density of states and Fermi statistics
- Expected lifetimes  $\sim 10-100\text{ms}$



G. M. Bruun and P. Massignan, Phys. Rev. Lett. 105, 020403 (2010).







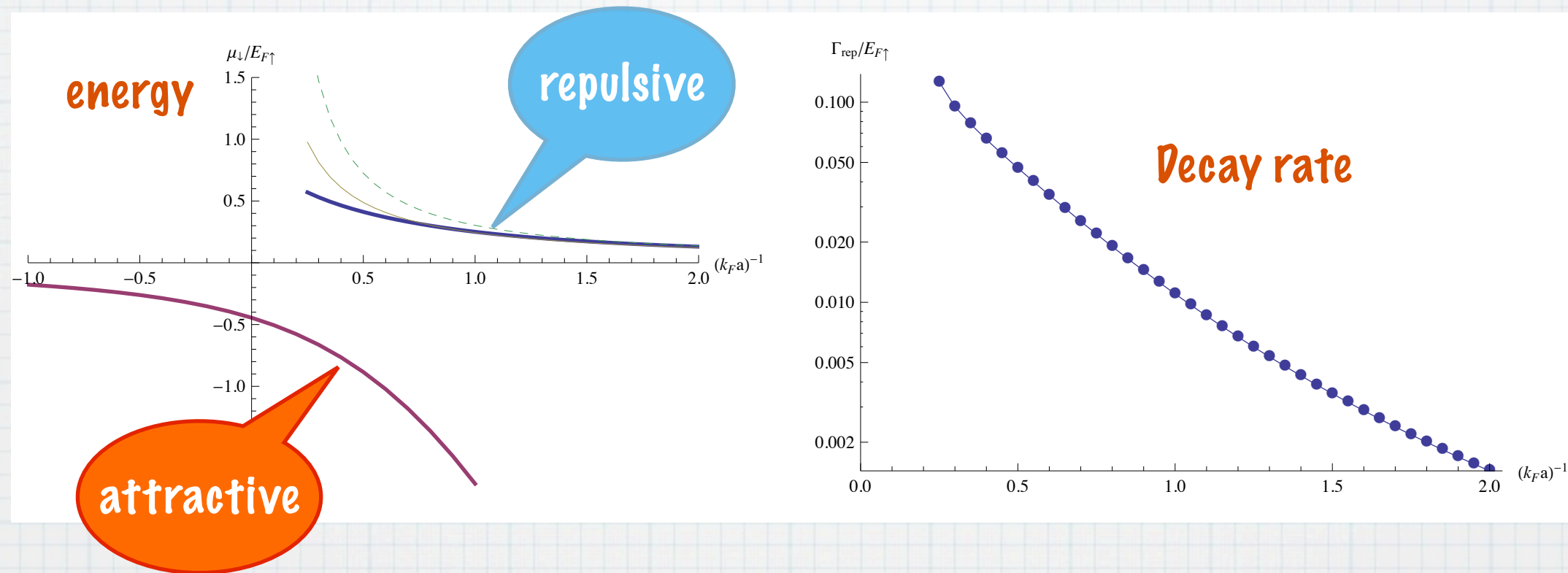
atom-molecule  
coupling:

$$\bullet = \frac{1}{g(\mathbf{p}, z)} = \int \frac{d^3 q}{(2\pi)^3} \phi_q \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} \sim -\sqrt{\frac{m_r^2 a}{2\pi}}$$

**Vacuum:**  $D_0(\mathbf{p}, z) = \int d^3 \check{q} \phi_q^2 \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} + \frac{T_2(\mathbf{p}, z)}{g(\mathbf{p}, z)^2}$



# Repulsive polaron



A  $^{40}\text{K}$  impurity in a Fermi sea of  $^6\text{Li}$



# Impurity spectral function

$$A = -\text{Im}[G_{\uparrow}(k=0, \omega)]$$

