

Impurities in a Fermi sea

Pietro Massignan (UAB&ICFO-Barcelona)

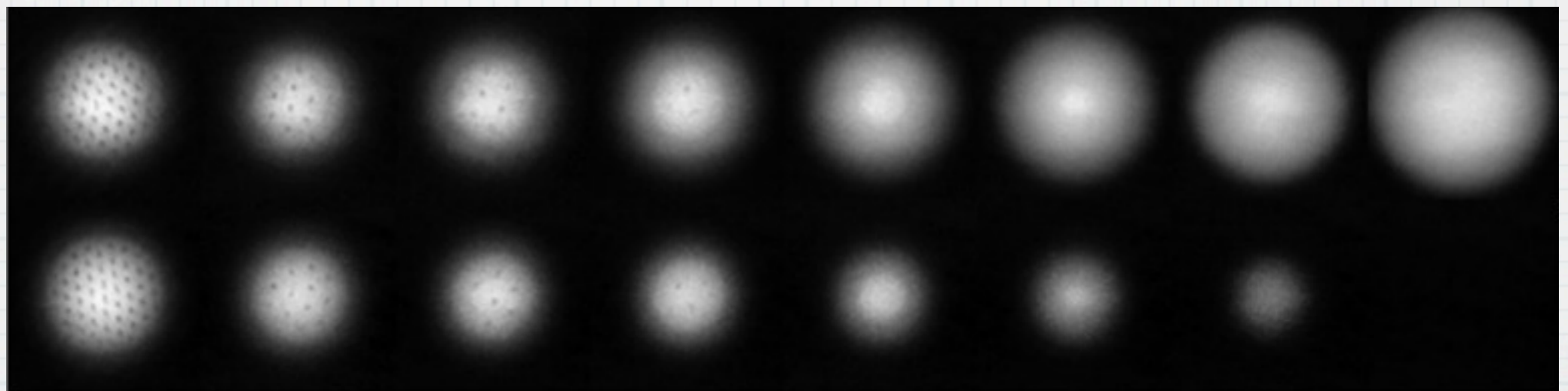


Fermi mixtures

BEC-BCS



SF-normal



$N=N$

2

$N \gg N$

in collaboration with:



Georg Bruun (Aarhus)



Carlos Lobo
(Southampton)



Alessio Recati
(Trento)



Kayvan Sadegzadeh (Cambridge)

Outline

- * Quasi-particles in many-body systems
- * The MIT “impurity” experiment
- * Polarons and molecules
- * Decay rates
- * Repulsive polaron

Many-body systems

2 A GUIDE TO FEYNMAN DIAGRAMS [0.0

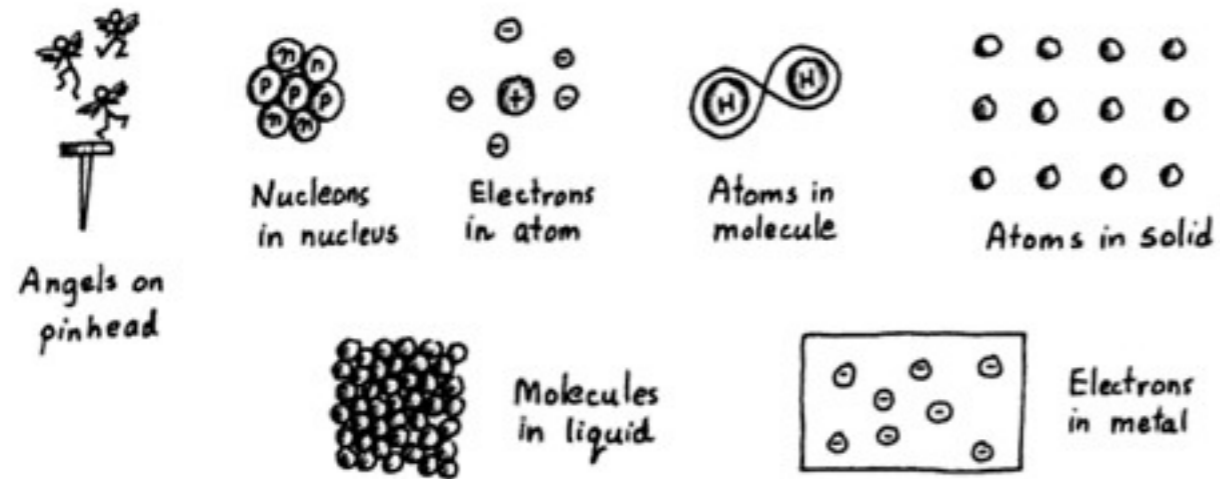
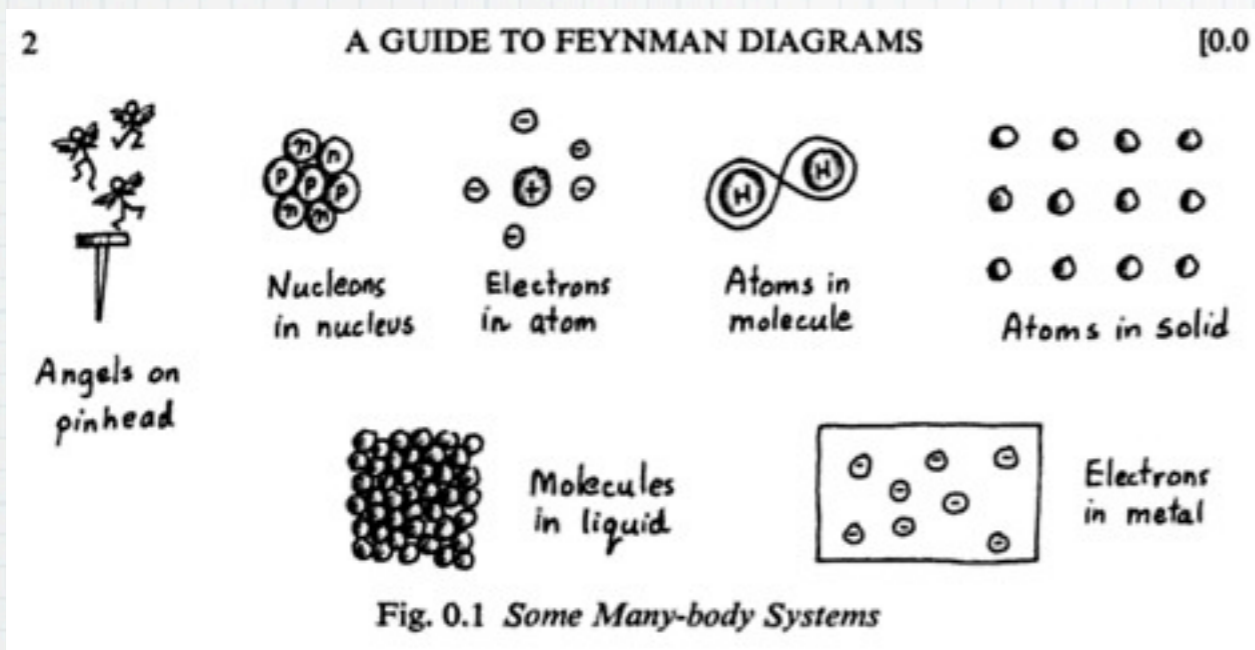


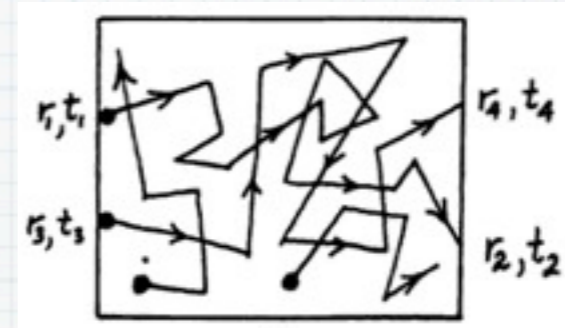
Fig. 0.1 *Some Many-body Systems*

(from Richard Mattuck's book)

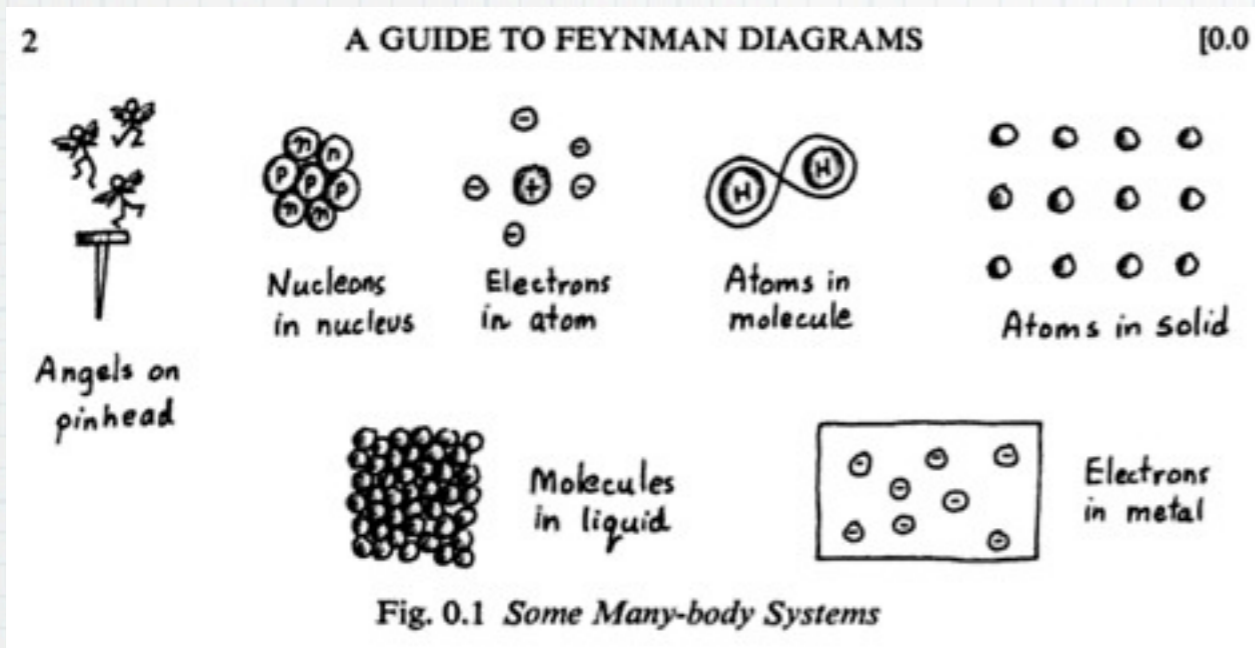
Many-body systems



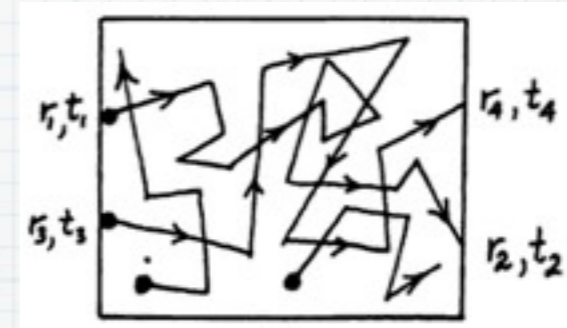
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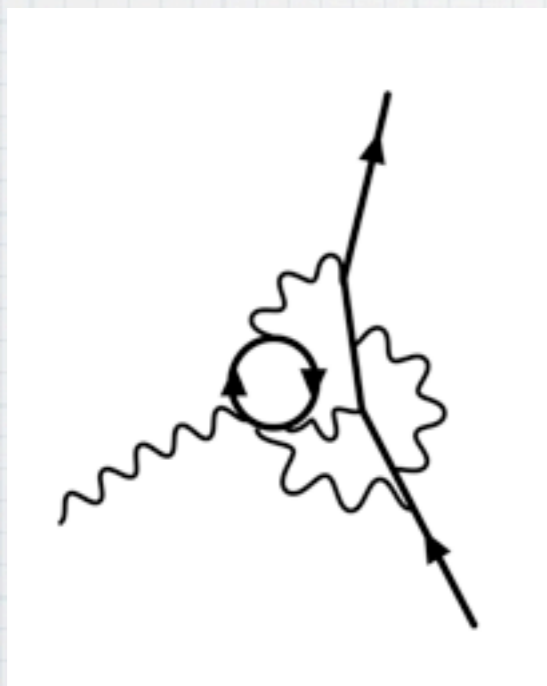
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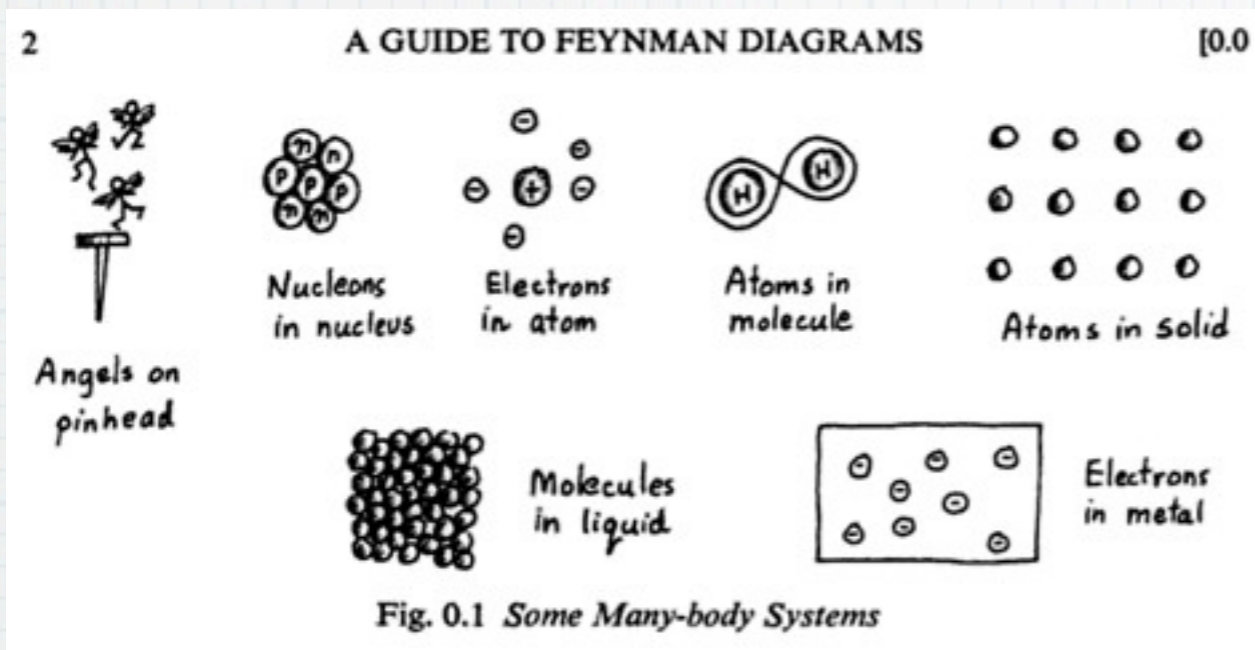
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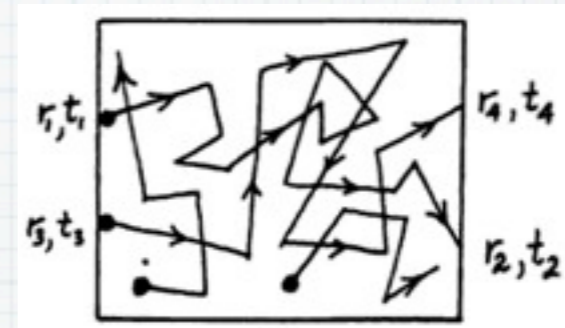
Feynman diagrams:



Many-body systems



(from Richard Mattuck's book)



Feynman diagrams:

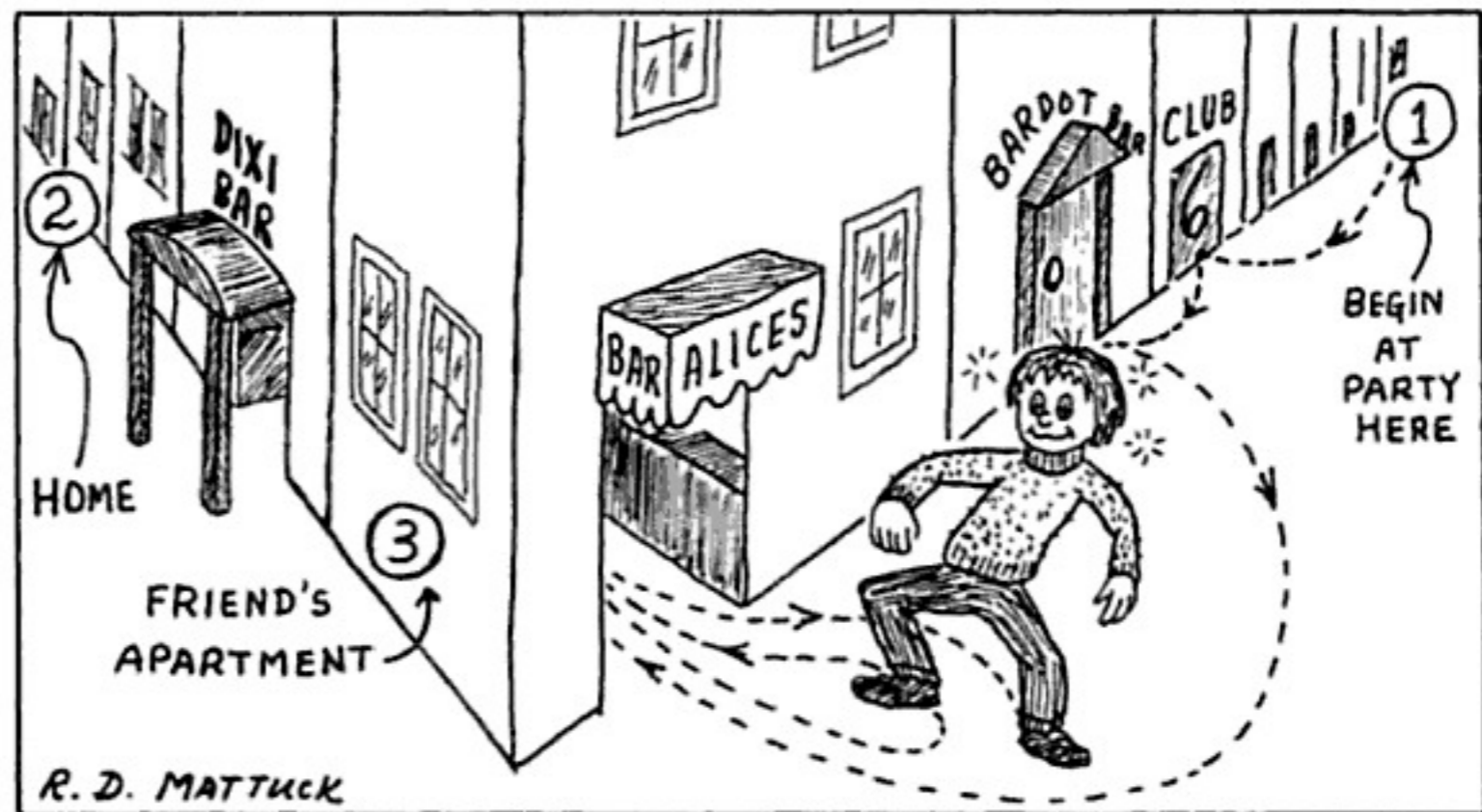
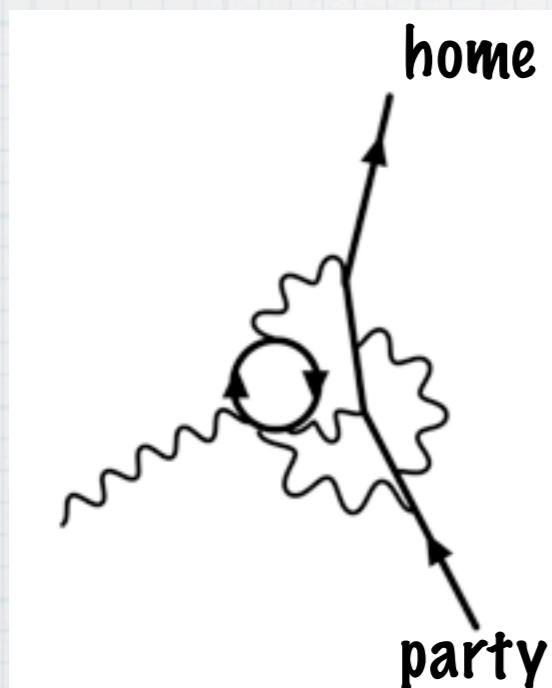
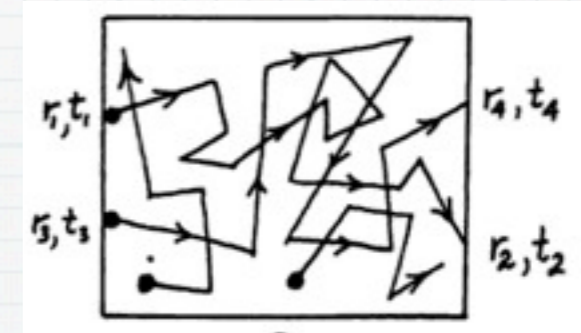


Fig. 1.1 Propagation of Drunken Man

Quasi-Particles

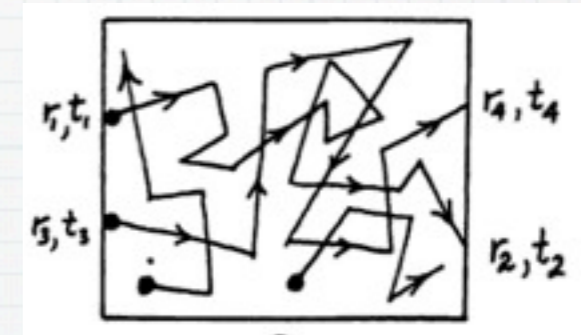
Landau's idea:
who cares about real particles?



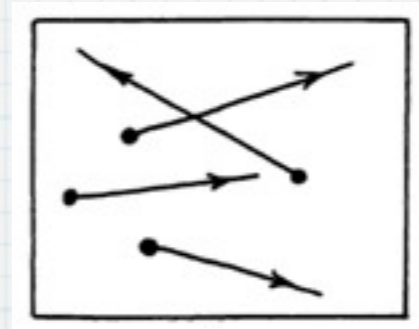
Of importance are the excitations,
which behave
as **quasi**-particles!

Quasi-Particles

Landau's idea:
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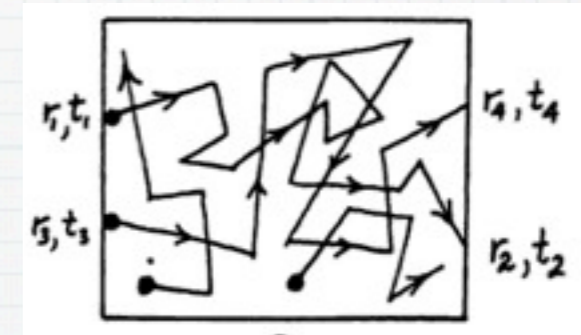


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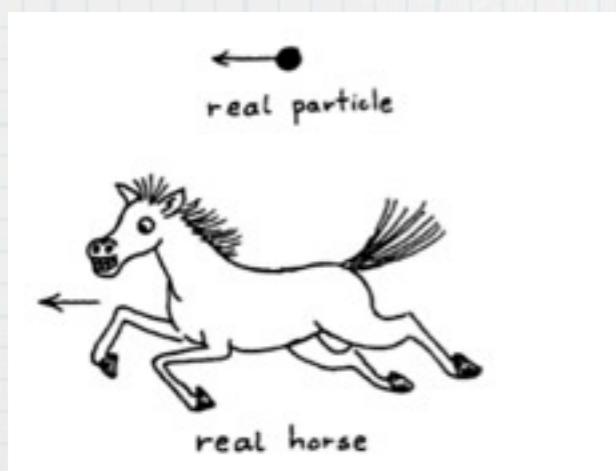
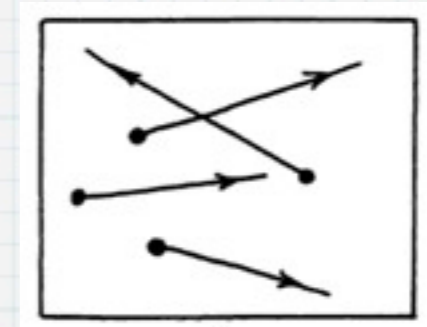


Quasi-Particles

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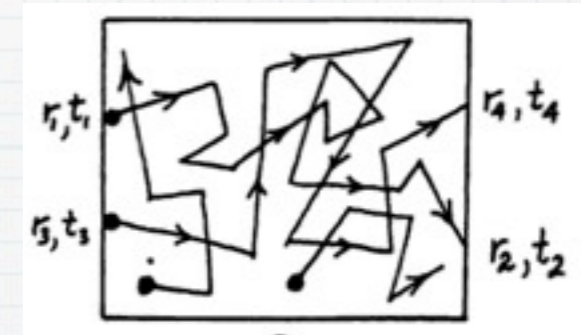


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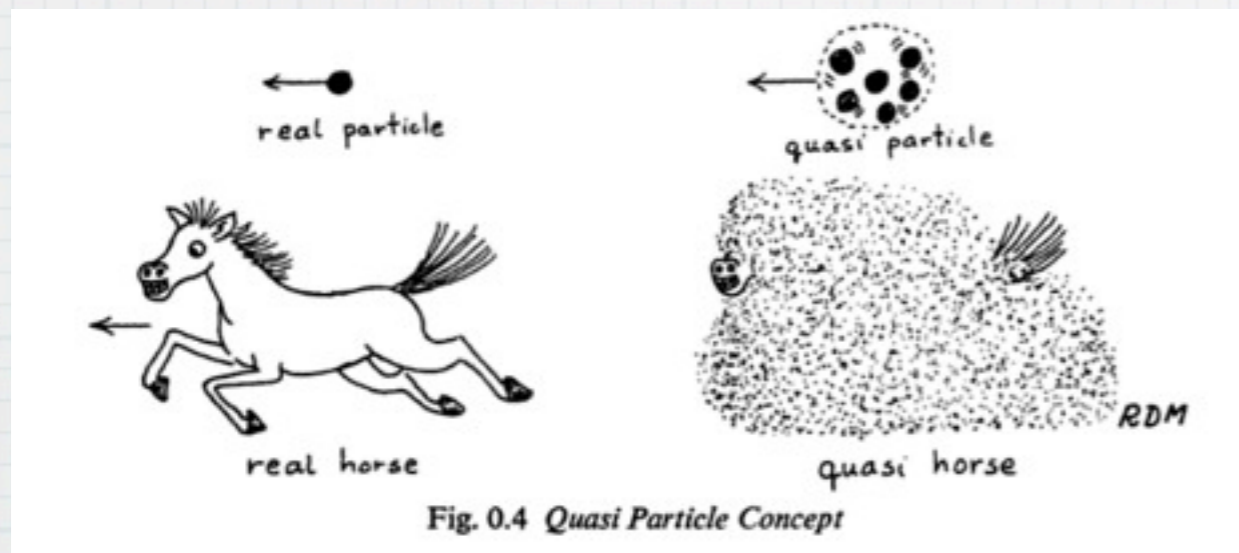
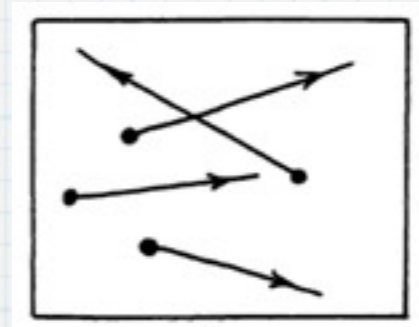


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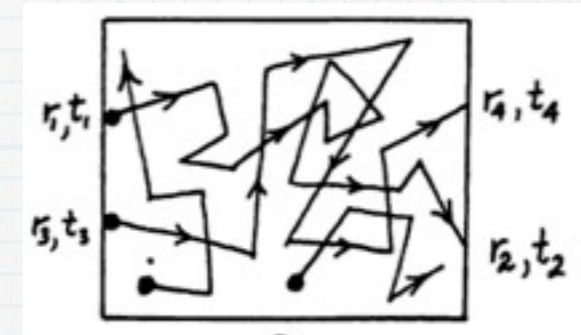


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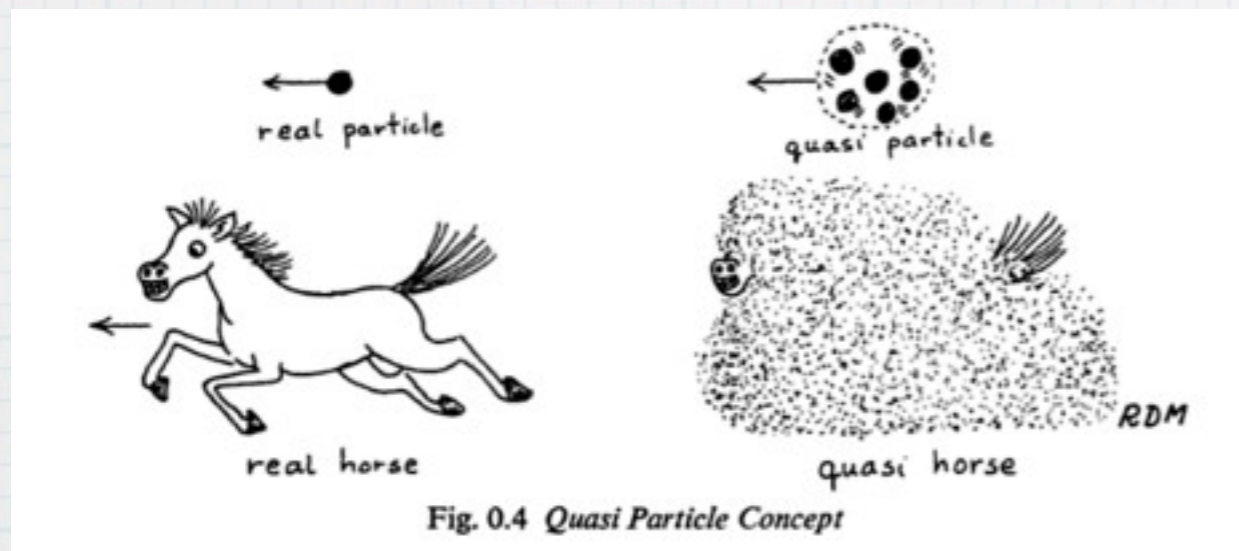
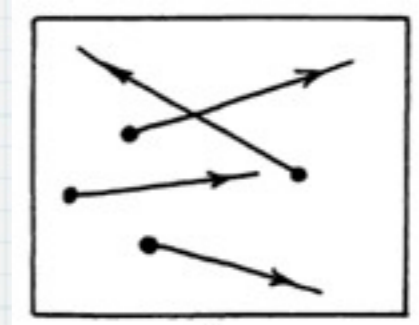


Quasi-Particles

Landau's idea:
who cares about real particles?



Of importance are the excitations,
which behave
as **quasi**-particles!



a **QP** is a "free particle" with:
@ **renormalized mass**
@ **chemical potential**
@ **shielded interactions**
@ **q. numbers (charge, spin, ...)**
@ **lifetime**

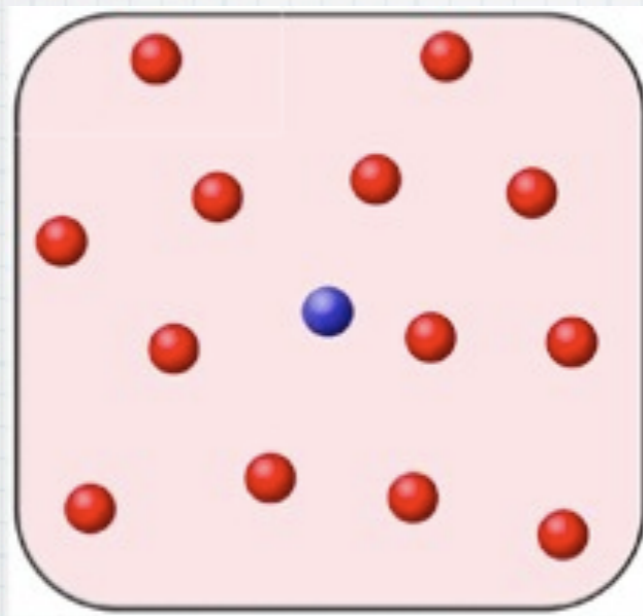
The MIT experiment

Schirotzek, Wu, Sommer & Zwierlein, PRL 2009

- non-interacting ↑ Fermi sea ($N \gg 1$)
- a single ↓ impurity

BCS $\xrightarrow{\text{Attraction strength}}$ BEC

$$(k_{Fa})^{-1} < 0$$



free particle

$$(k_{Fa})^{-1} > 0$$

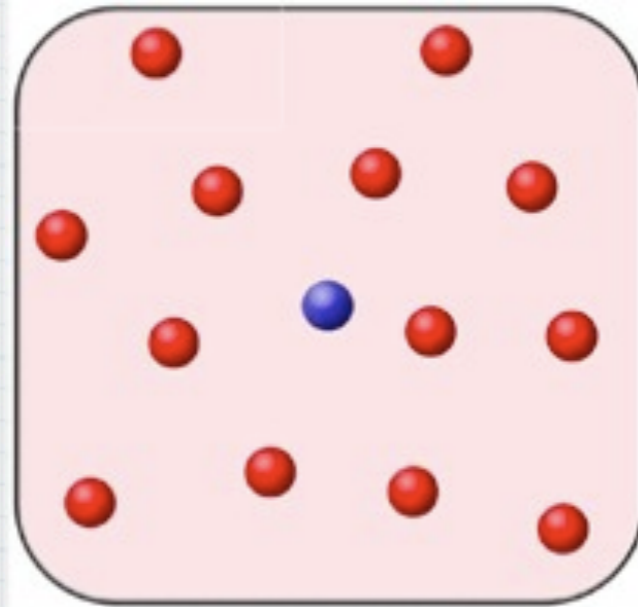
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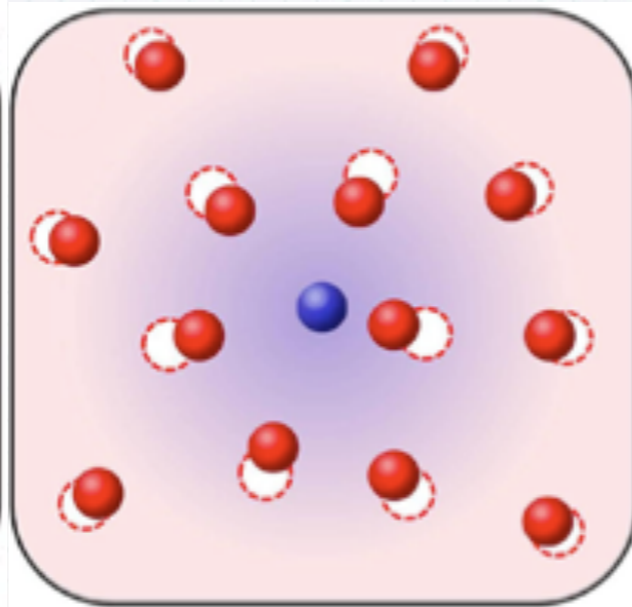
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QP (polaron)

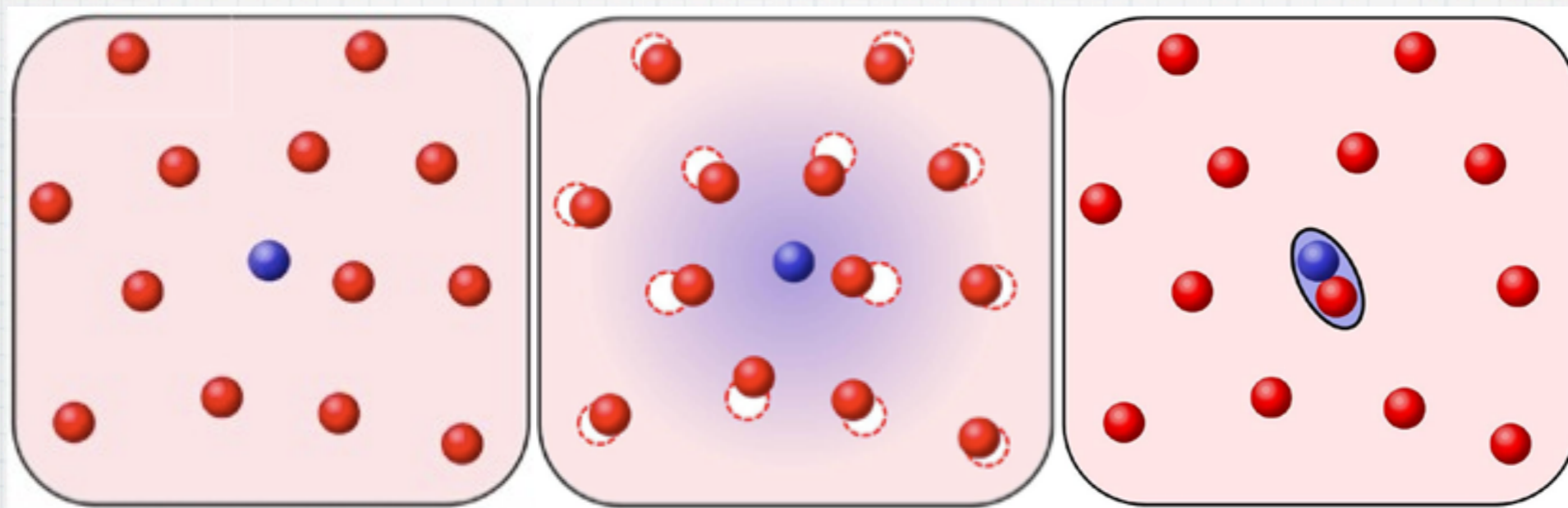
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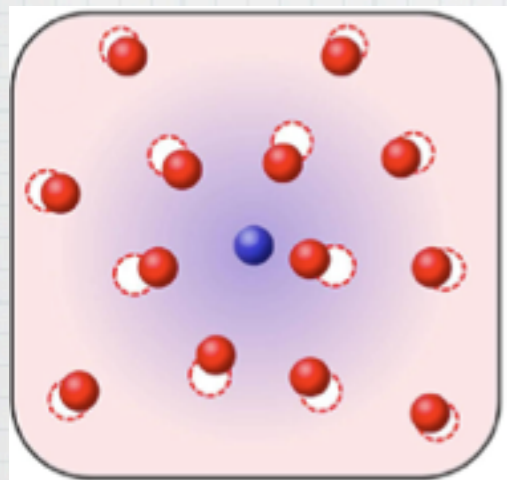
QP (polaron)

QP (molecule)

P-M transition: Prokof'ev & Svistunov, PRB 2008

Polaron: variational Ansatz

(F. Chevy, PRA 2006)



Polaron: variational Ansatz

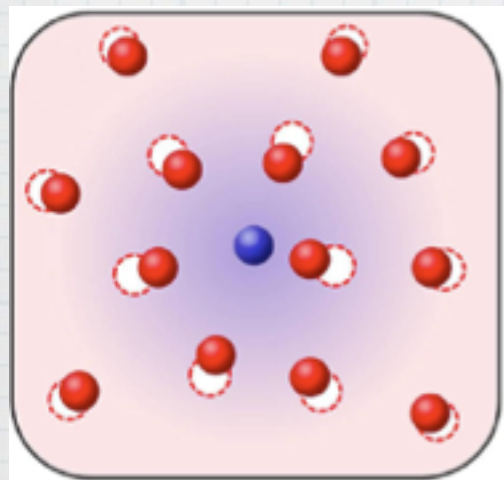
(F. Chevy, PRA 2006)

the ↓ impurity

$$|\psi_{\mathbf{p}}\rangle = \phi_0 c_{\mathbf{p}\downarrow}^\dagger |0\rangle_\uparrow + \sum_{q < k_F} \phi_{\mathbf{q}\mathbf{k}} c_{\mathbf{p}+\mathbf{q}-\mathbf{k}\downarrow}^\dagger c_{\mathbf{k}\uparrow}^\dagger c_{\mathbf{q}\uparrow} |0\rangle_\uparrow$$

non-interacting ↑ Fermi sea

Particle-Hole dressing



Polaron: variational Ansatz

(F. Chevy, PRA 2006)

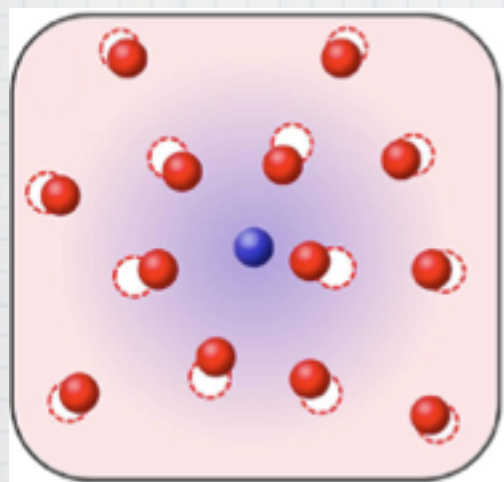
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Particle-Hole dressing

Very good agreement with QMC results for μ_\downarrow and m^*



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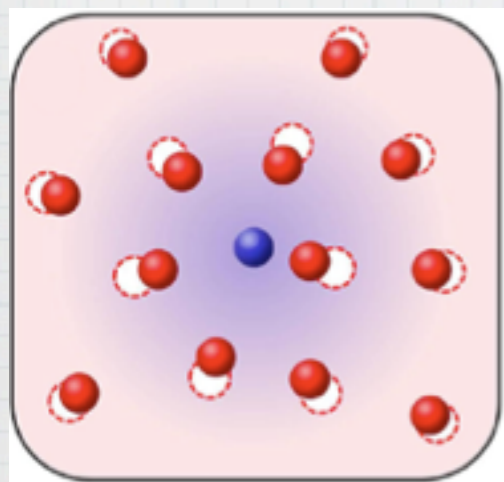
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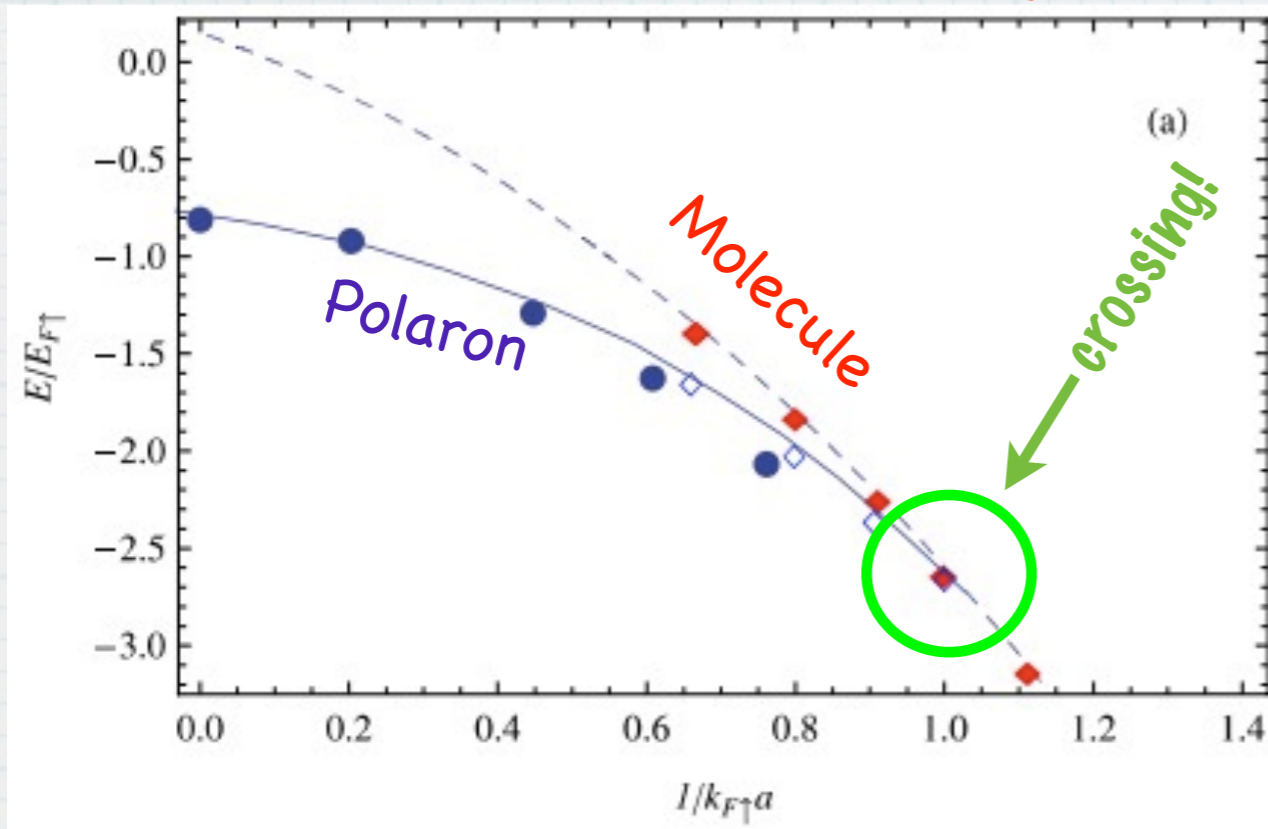
Very good agreement with QMC results for μ_\downarrow and m^*

This variational Ansatz has a diagrammatic equivalent: the forward scattering, or ladder, approximation.

(Combescot et al., PRL 2007)

QP parameters

Chemical potential $\mu \downarrow$



◇, ◆ : QMC
—, - - : variat, diagr
● : MIT expmt

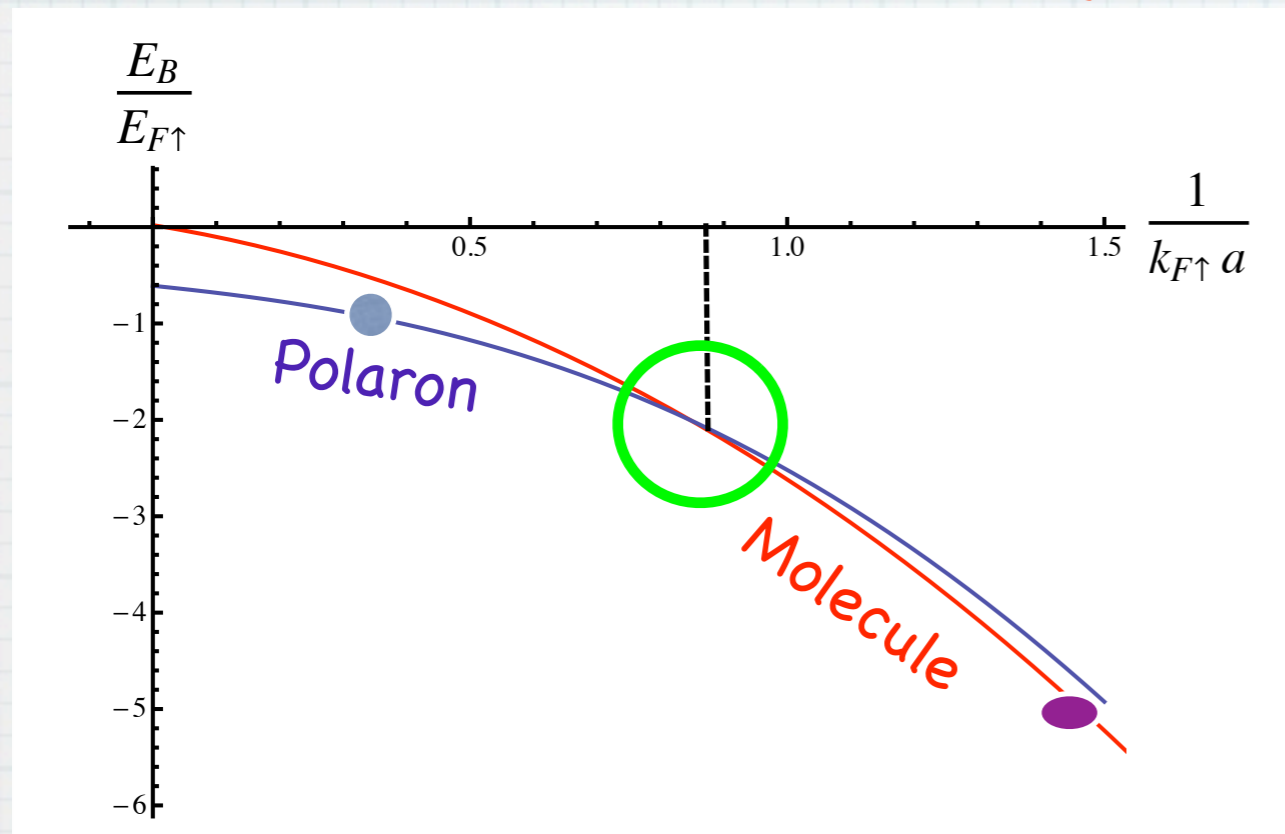
QMC: Prokof'ev&Svistunov

Variational and diagrammatic: Chevy, Recati, Lobo, Stringari, Combescot, Leyronas
Massignan&Bruun, Zwirger, Punk, Stoof, Mora,...

Experiments: MIT, ENS

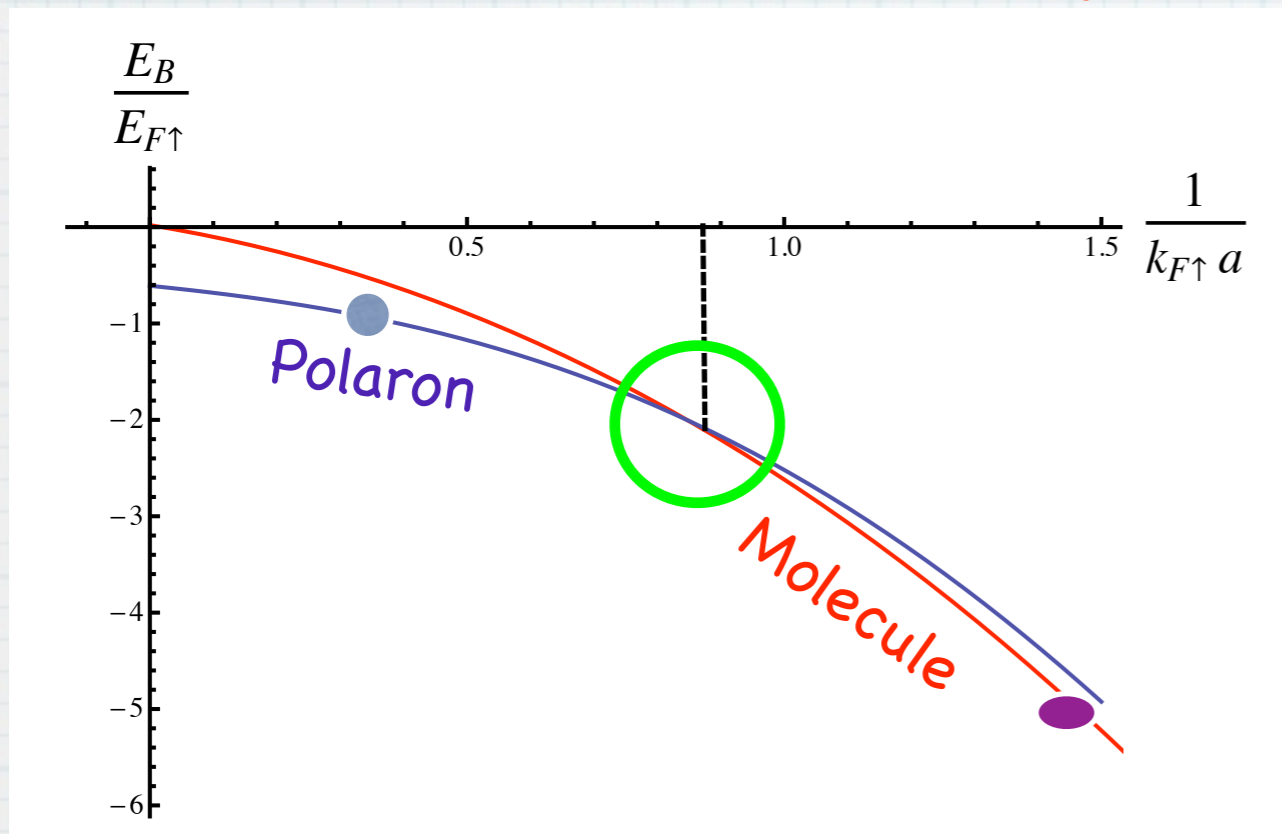
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Chemical potential μ_{\downarrow}

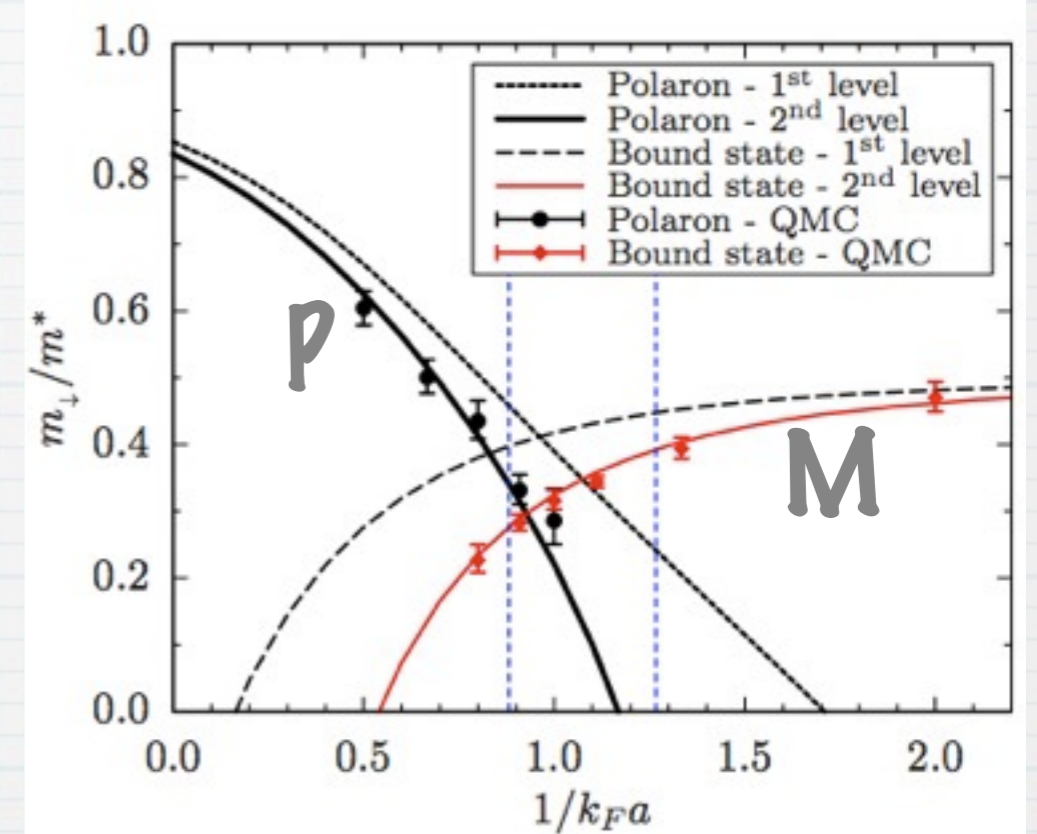


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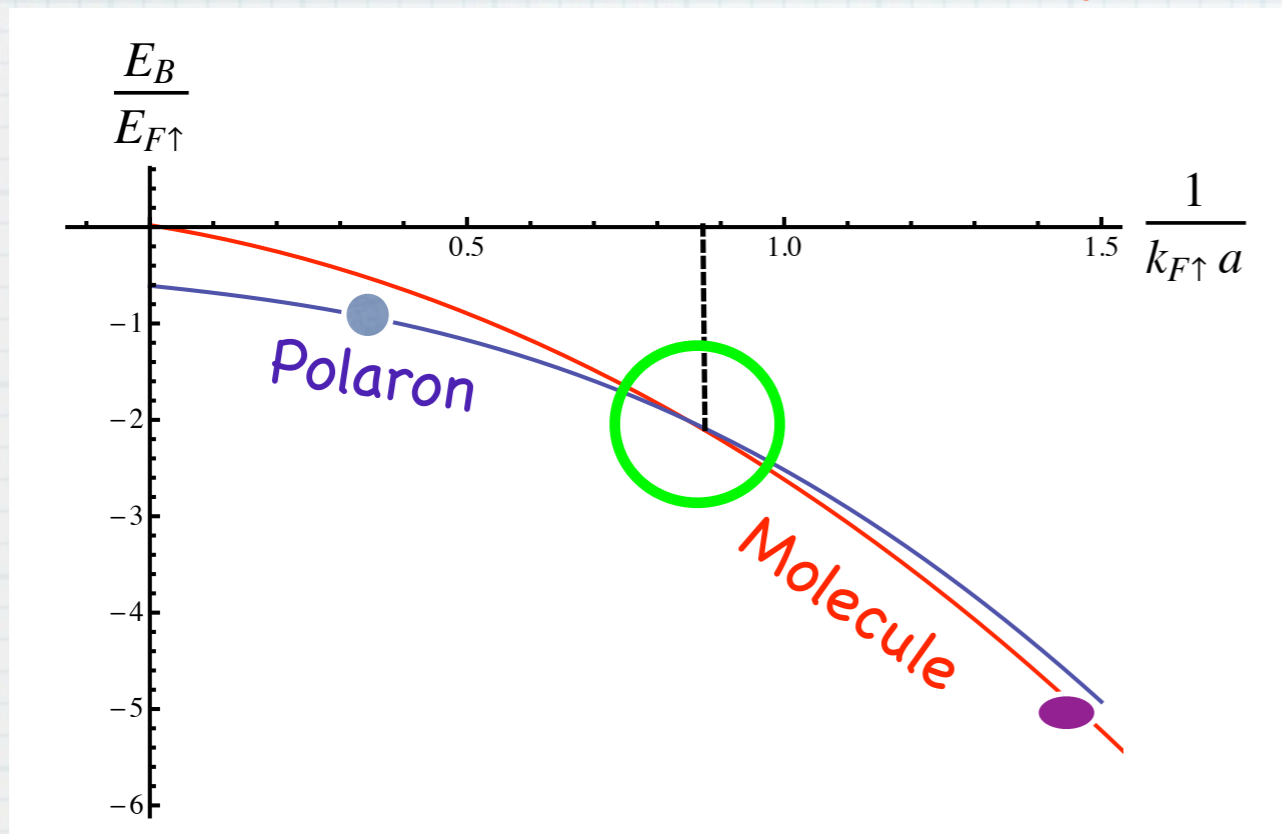


Effective mass m^*

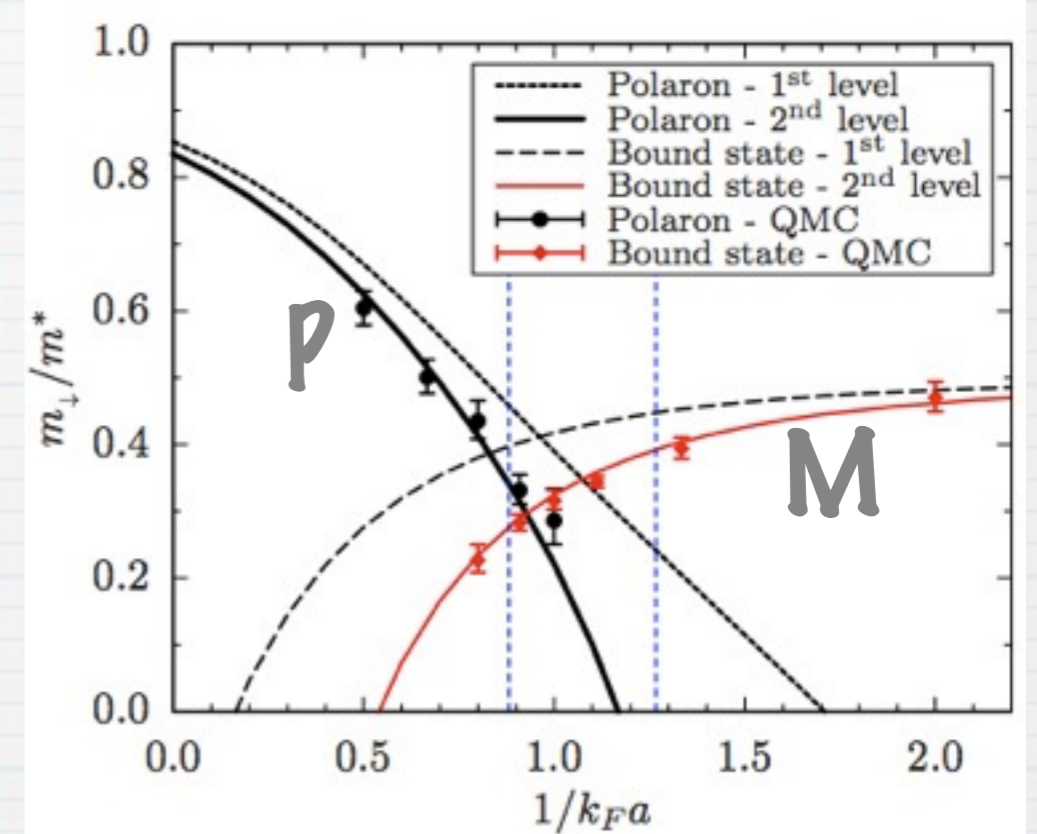


QP parameters

Chemical potential μ_{\downarrow}



Effective mass m^*



P-P Interactions: Mora&Chevy, PRL 2010

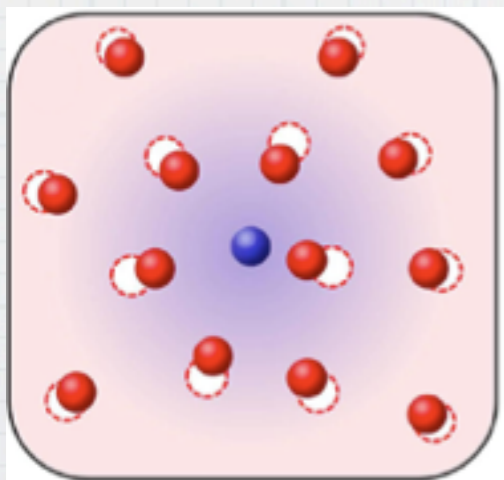
Zhenhua, Zöllner & Pethick, PRL 2010

Equation of state of a unitary Fermi gas

In the normal phase at $T=0$,

$$P = \overset{\text{non-interacting } \uparrow}{\frac{1}{15\pi^2} \left(\frac{2m_{\uparrow}}{\hbar^2}\right)^{3/2} \mu_{\uparrow}^{5/2}} + \overset{\text{non-interacting QP}}{\frac{1}{15\pi^2} \left(\frac{2m_{\downarrow}^*}{\hbar^2}\right)^{3/2} (\mu_{\downarrow} - A\mu_{\uparrow})^{5/2}}$$

$$A = -0.615$$
$$m_{\downarrow}^* = 1.2m$$



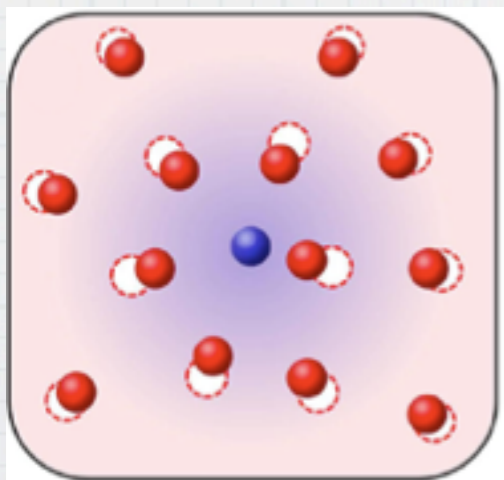
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$$A = -0.615$$

$$m_{\downarrow}^* = 1.2m$$



Same thermodynamics for:

- ultracold atoms
- dilute neutron matter

What's left?

What's left?

- ✓ chemical potential
- ✓ renormalized mass
- ✓ shielded interactions
- ✓ lifetime



G. Bruun & PM, PRL 2010

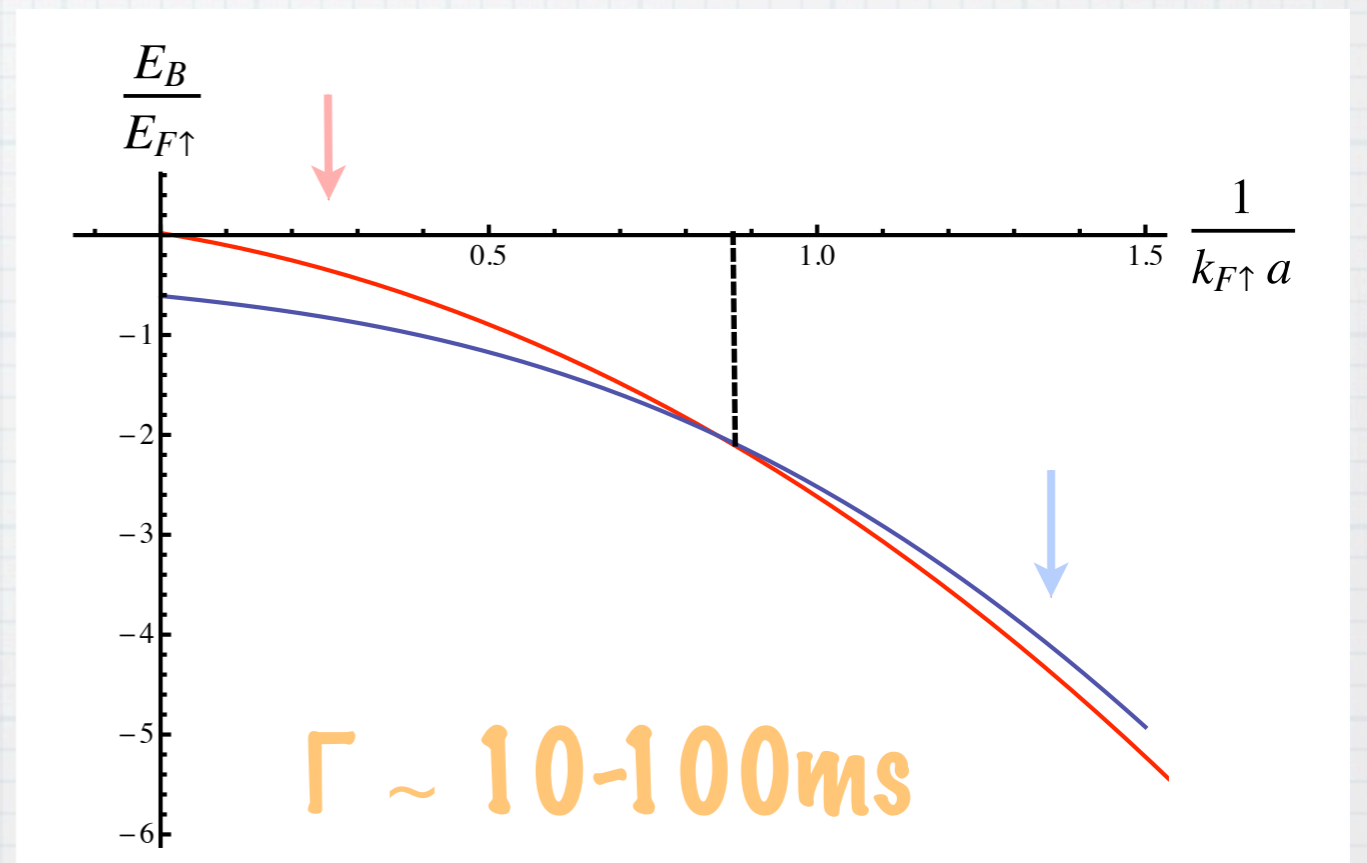
Very long QP lifetimes!

G. Bruun & PM, PRL 2010

$$\Gamma_P \sim Z_M (\Delta\omega)^{9/2}$$

$$\Delta\omega = \omega_P - \omega_M$$

$$\Gamma_M \sim Z_P (-\Delta\omega)^{9/2}$$



Pol \rightarrow Mol decay

$$\Delta\omega = \omega_P - \omega_M > 0$$

Polaron: $G_{\downarrow}(\mathbf{p}, z)^{-1} = G_{\downarrow}^0(\mathbf{p}, z)^{-1} - \Sigma_P(\mathbf{p}, z)$

Decay rate: $\Gamma_P = -\text{Im}\Sigma_P(p=0, \omega_P)$

Pol \rightarrow Mol decay

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Hole expansion: $\Sigma_P(\mathbf{p}, z) = \Sigma_P^{(1)}(\mathbf{p}, z) + \Sigma_P^{(2)}(\mathbf{p}, z) + \dots$

Pol \rightarrow Mol decay

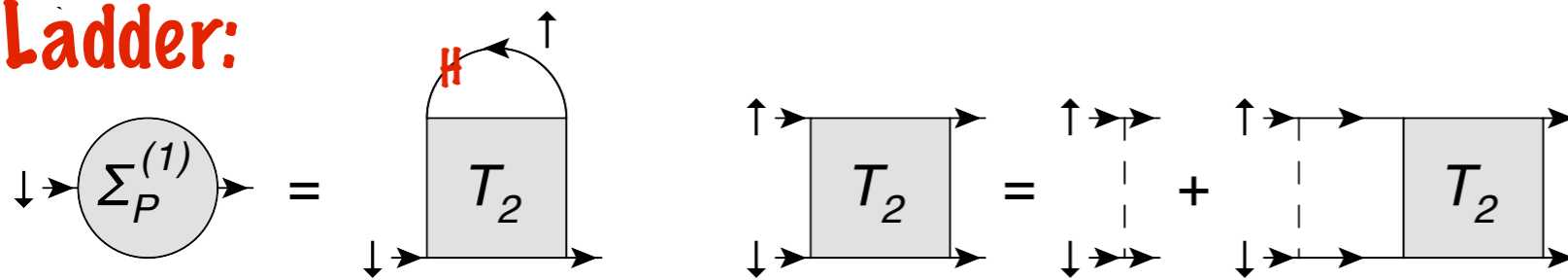
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Ladder:



**no damping
in the ladder approx.**

Pol \rightarrow Mol decay

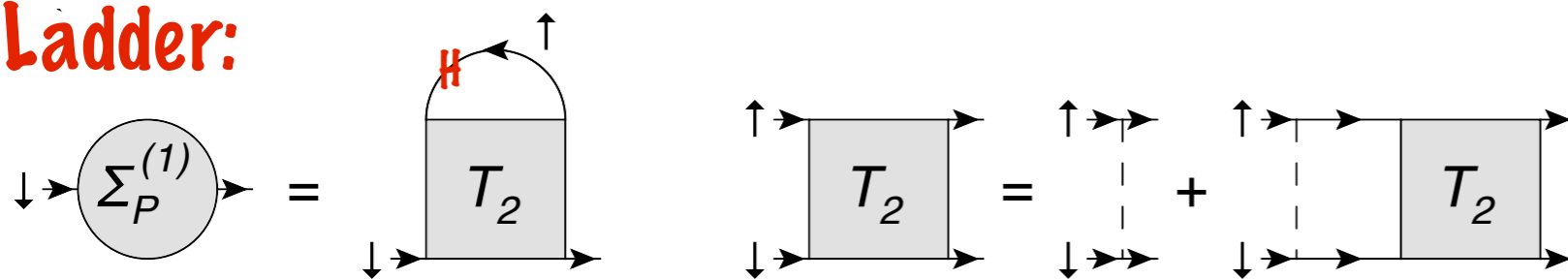
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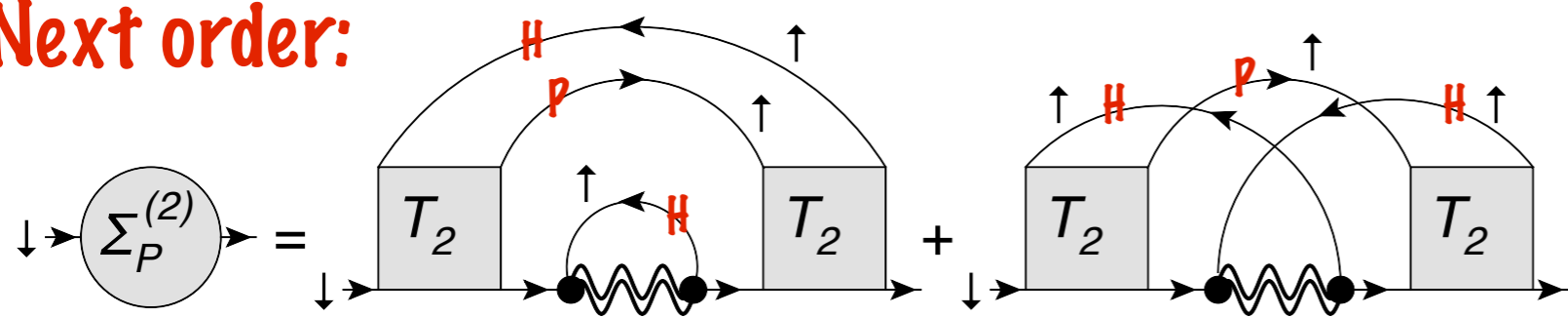
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Ladder:



**no damping
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Next order:



dressed molecule

Pol \rightarrow Mol decay

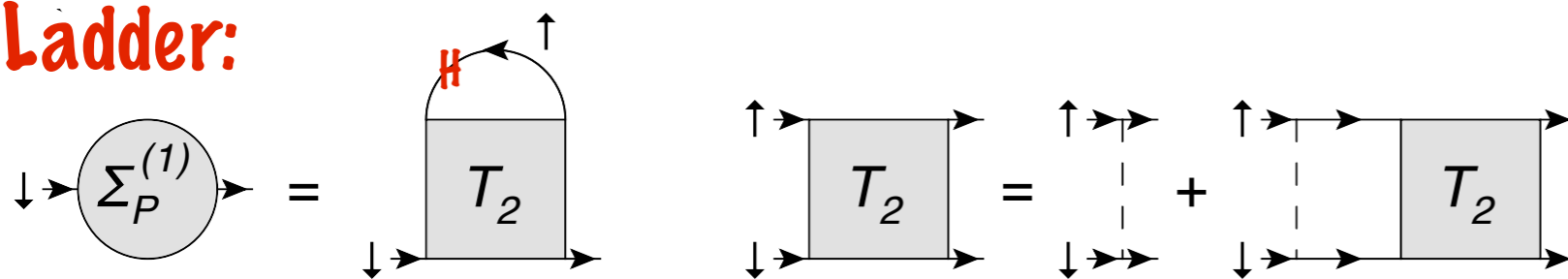
$$\Delta\omega = \omega_P - \omega_M > 0$$

Polaron: $G_{\downarrow}(\mathbf{p}, z)^{-1} = G_{\downarrow}^0(\mathbf{p}, z)^{-1} - \Sigma_P(\mathbf{p}, z)$

Decay rate: $\Gamma_P = -\text{Im}\Sigma_P(p=0, \omega_P)$

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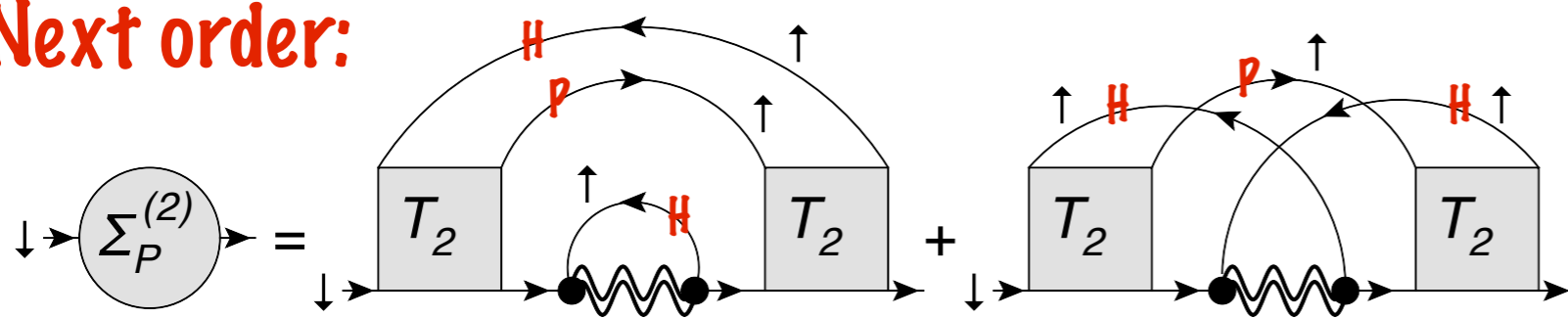
Ladder:



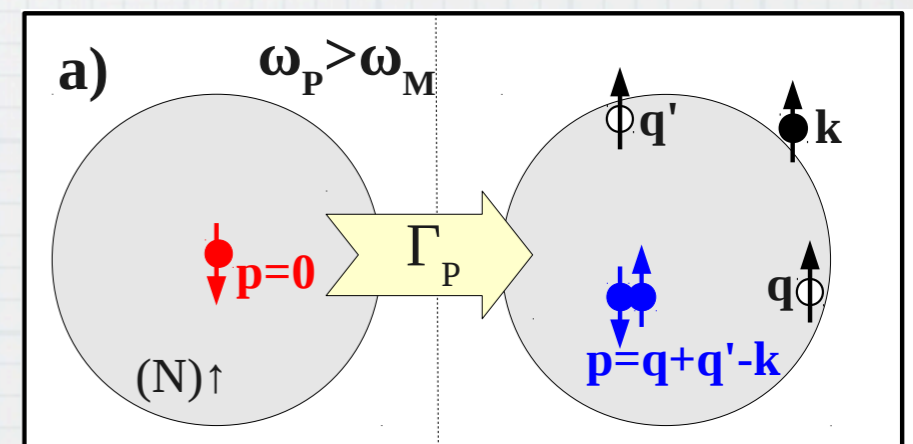
no damping
in the ladder approx.

3-body process

Next order:



dressed molecule



Fermi's Golden rule

atom-molecule coupling



matrix element $\sim \Delta\omega$



density of final states $\sim \Delta\omega^{1/2}$



Fermi's Golden rule

atom-molecule coupling

$$\Gamma_P = \frac{g^2 Z_M}{2} \int \frac{d^3 k d^3 q d^3 q'}{(2\pi)^9} [F(\mathbf{q}, \mathbf{k}, \omega_P) - F(\mathbf{q}', \mathbf{k}, \omega_P)]^2 \delta \left(\Delta\omega + \xi_{q\uparrow} + \xi_{q'\uparrow} - \xi_{k\uparrow} - \frac{(\mathbf{q} + \mathbf{q}' - \mathbf{k})^2}{2m_M^*} \right)$$

matrix element $\sim \Delta\omega$

density of final states $\sim \Delta\omega^{7/2}$

$$q, q' < k_F, \quad k > k_F$$

$$F(\mathbf{q}, \mathbf{k}, \omega) = T_2(\mathbf{q}, \omega + \xi_{q\uparrow}) G_{\downarrow}^0(\mathbf{q} - \mathbf{k}, \omega + \xi_{q\uparrow} - \xi_{k\uparrow})$$

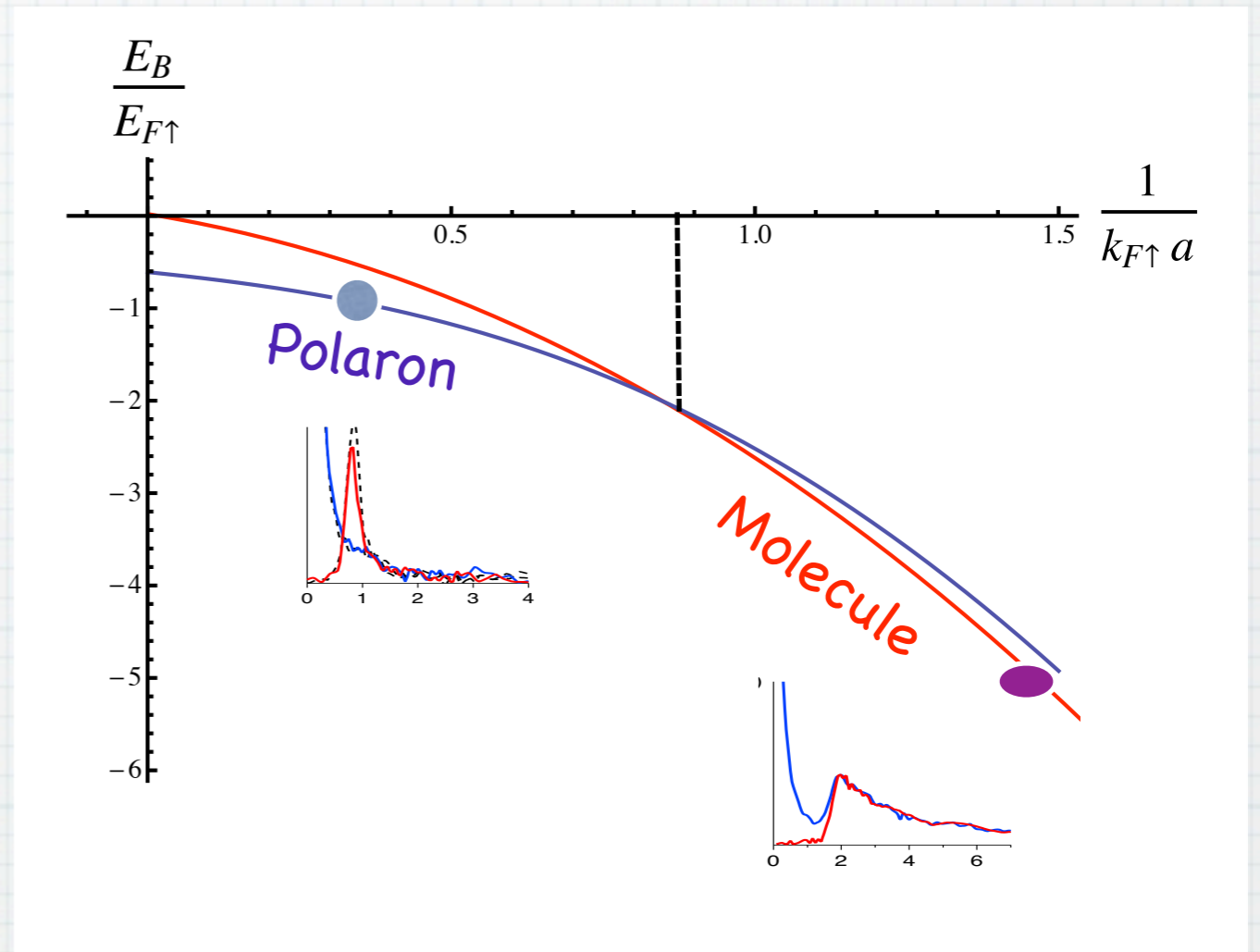
Experimental observation

Methods:

- RF spectra
- Collective modes to measure m^* vs. time
- Density profiles in the trap

Issues:

- * Phase separation?
 - * stabilized by finite T
 - * work with $m_{\downarrow} \neq m_{\uparrow}$
 - * use bosonic impurities
- * No decay to deeply bound molecular states



...is there more?

Yes..

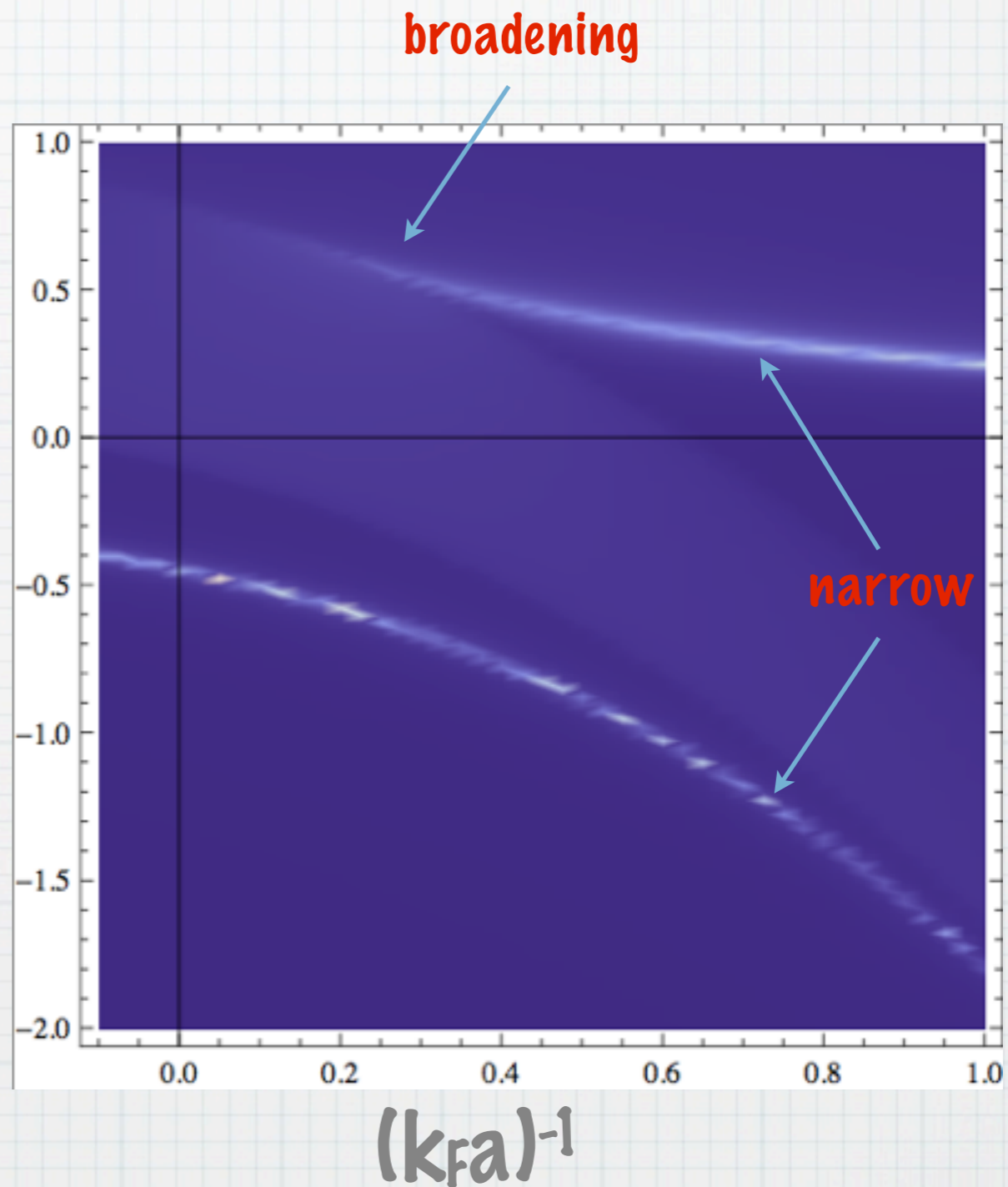
...is there more?

Yes..

spectral function

$$A_{\downarrow}(\omega) = -\text{Im}[G_{\downarrow}(k=0, \omega + i0^+)]$$

energy



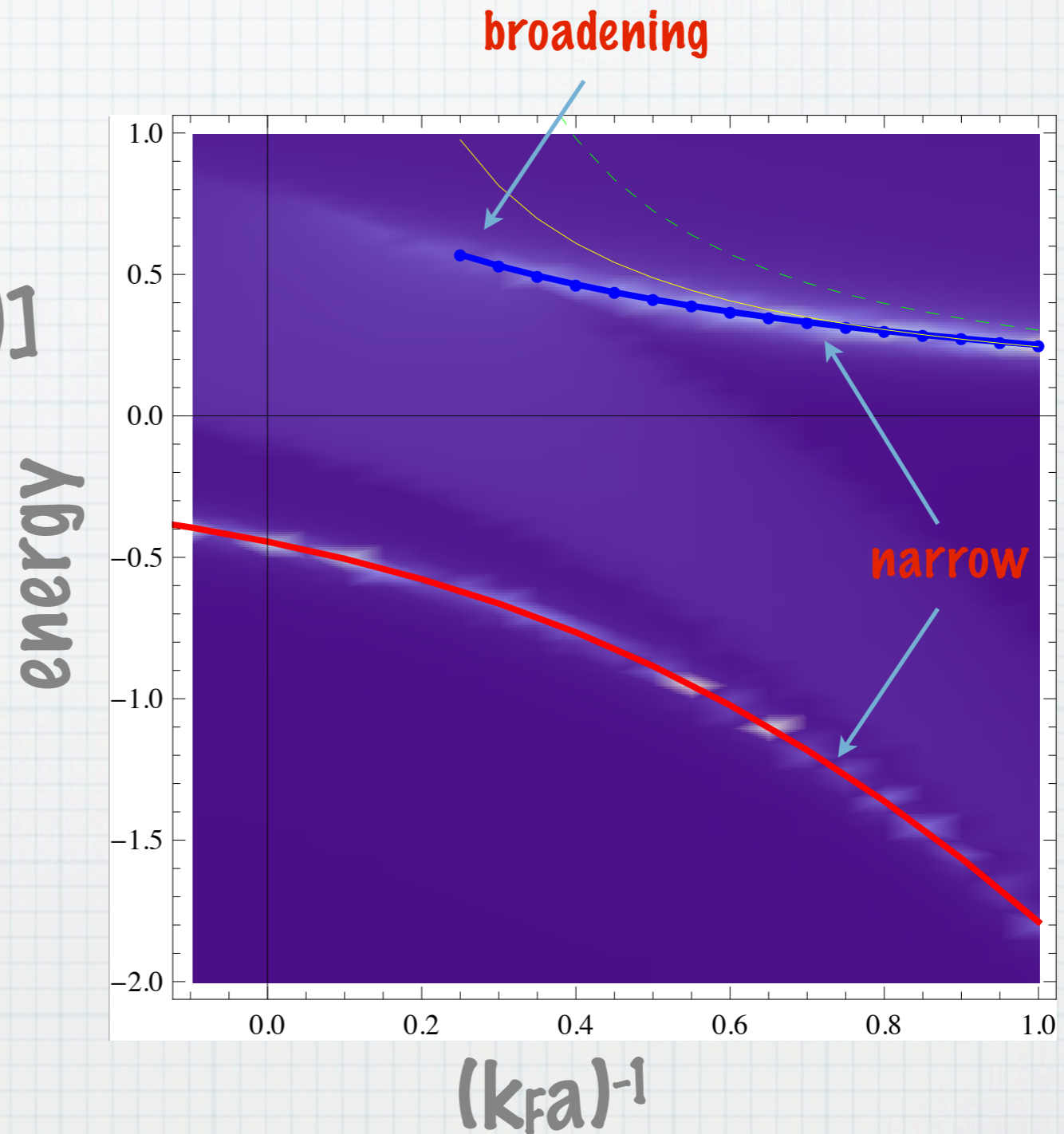
^{40}K impurity in a Fermi sea of ^6Li

...is there more?

Yes..

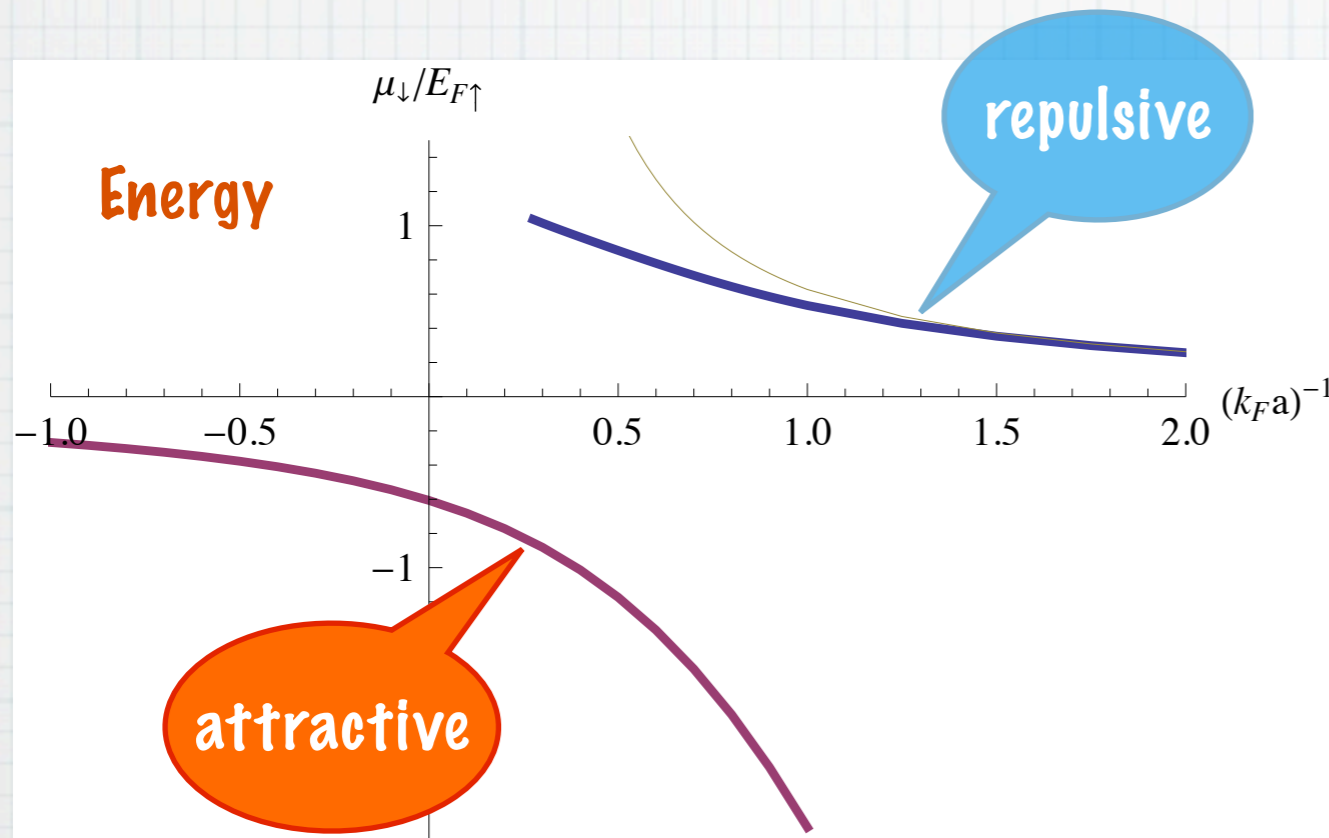
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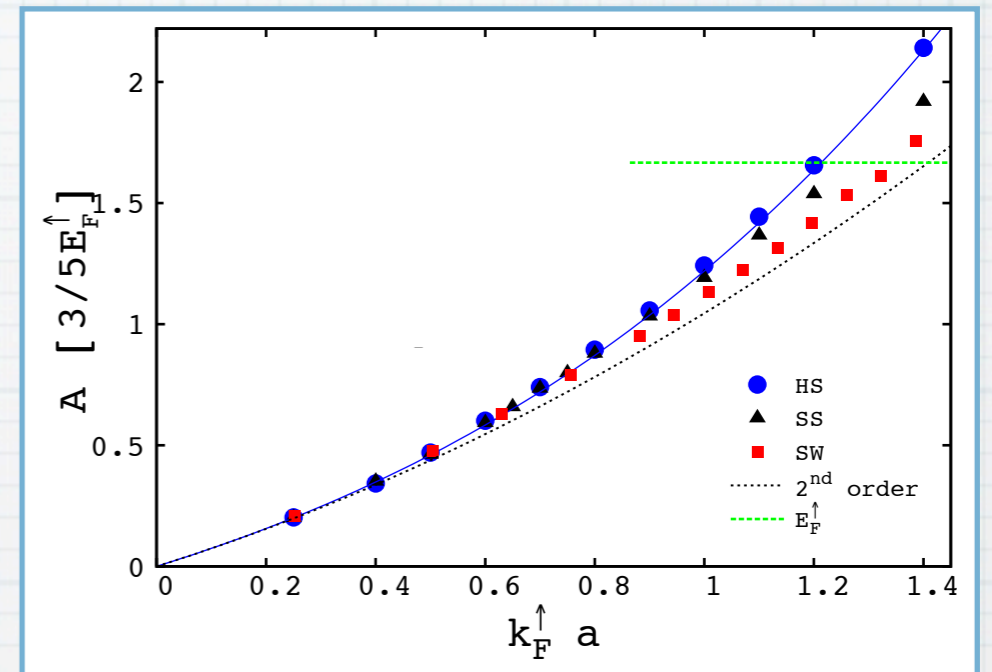


^{40}K impurity in a Fermi sea of ^6Li

Polaron energies



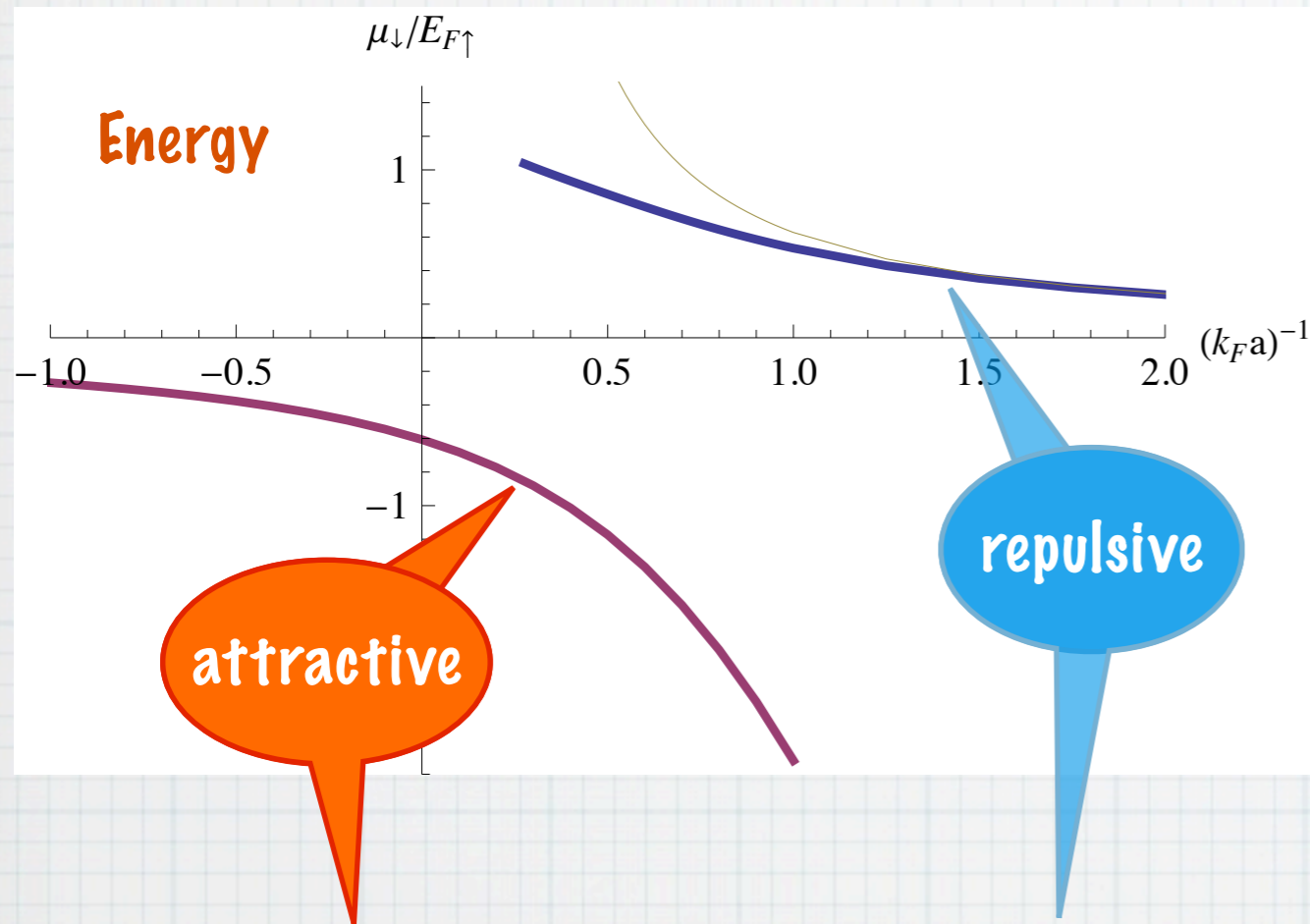
Massignan & Bruun
(in preparation)



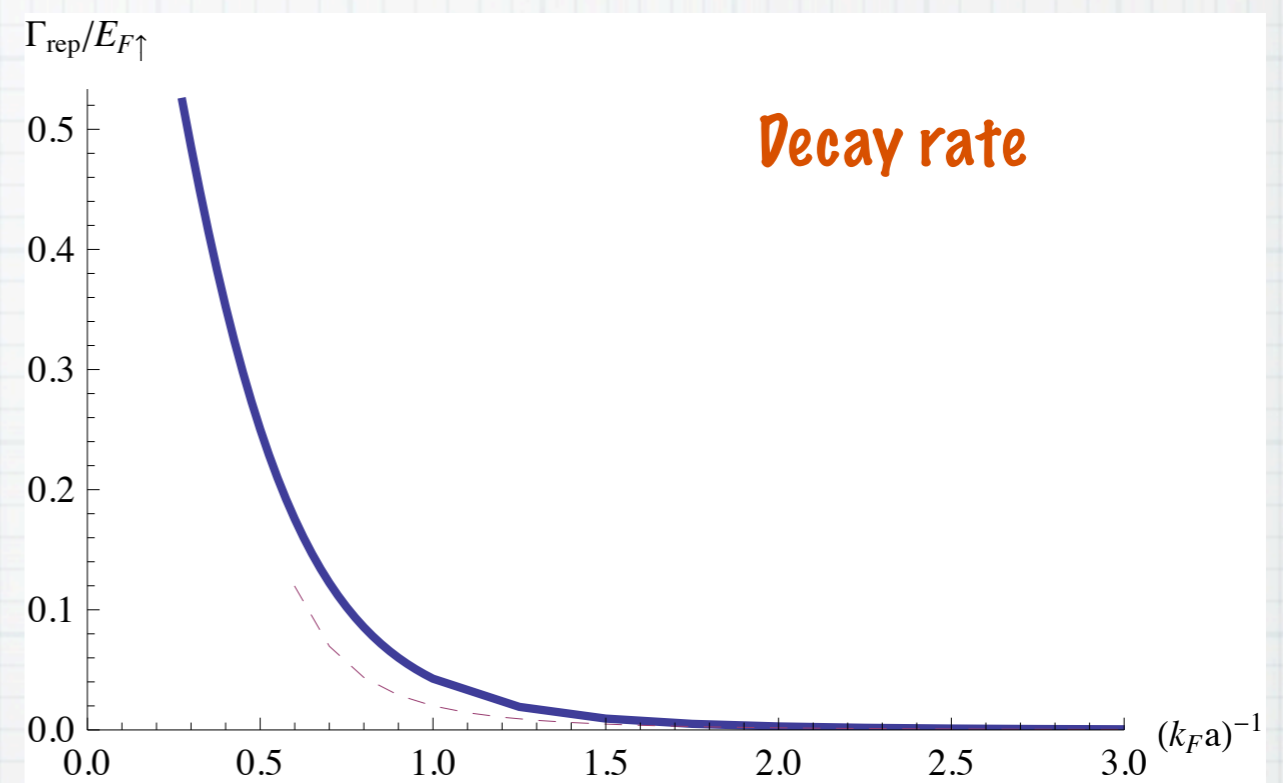
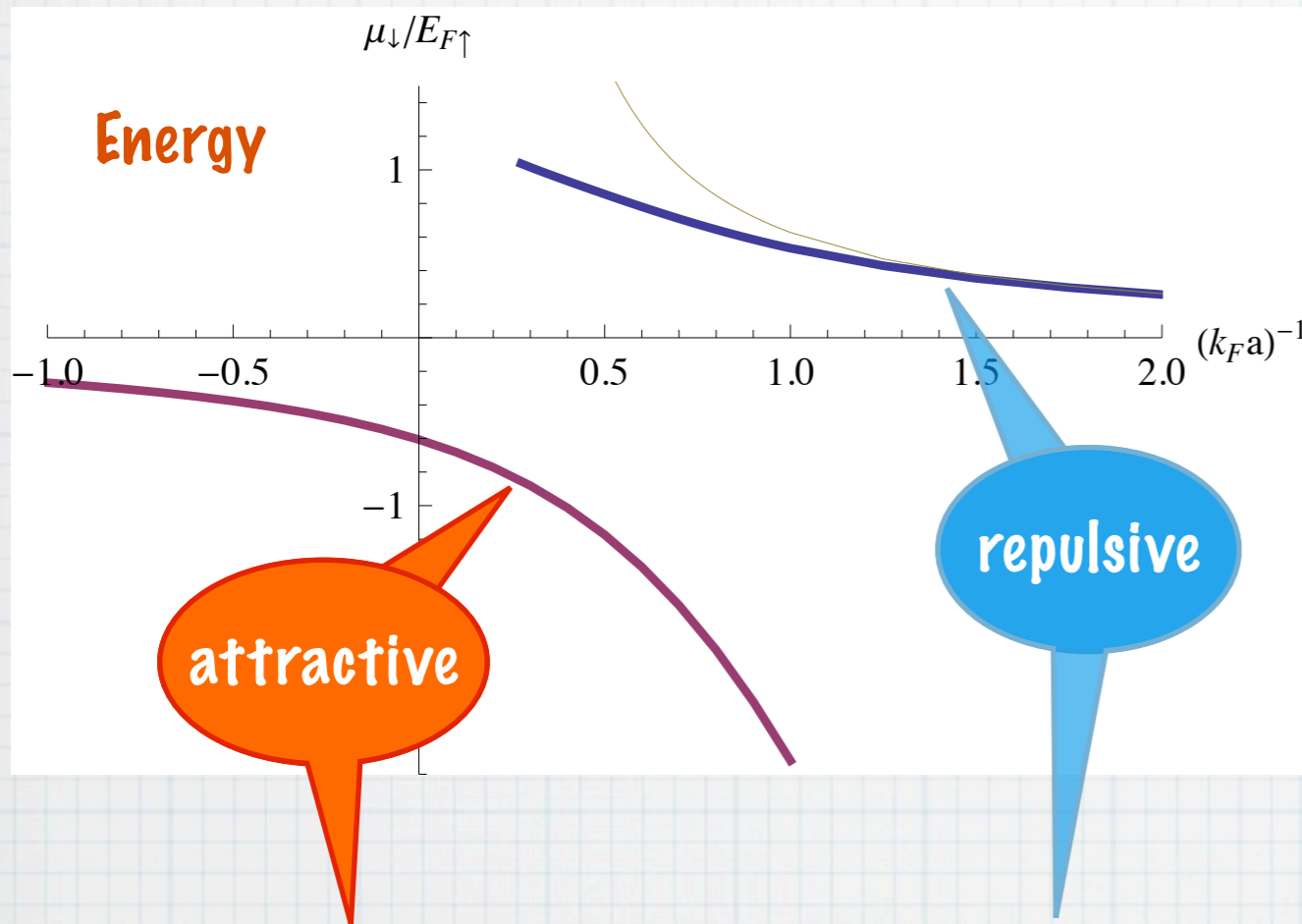
QMC by Pilati et al., PRL 2010

weak coupling:
$$\frac{\mu_\downarrow}{E_{F\uparrow}} = \frac{4}{3\pi} (k_F a) + \frac{2}{\pi^2} (k_F a)^2 + \dots$$

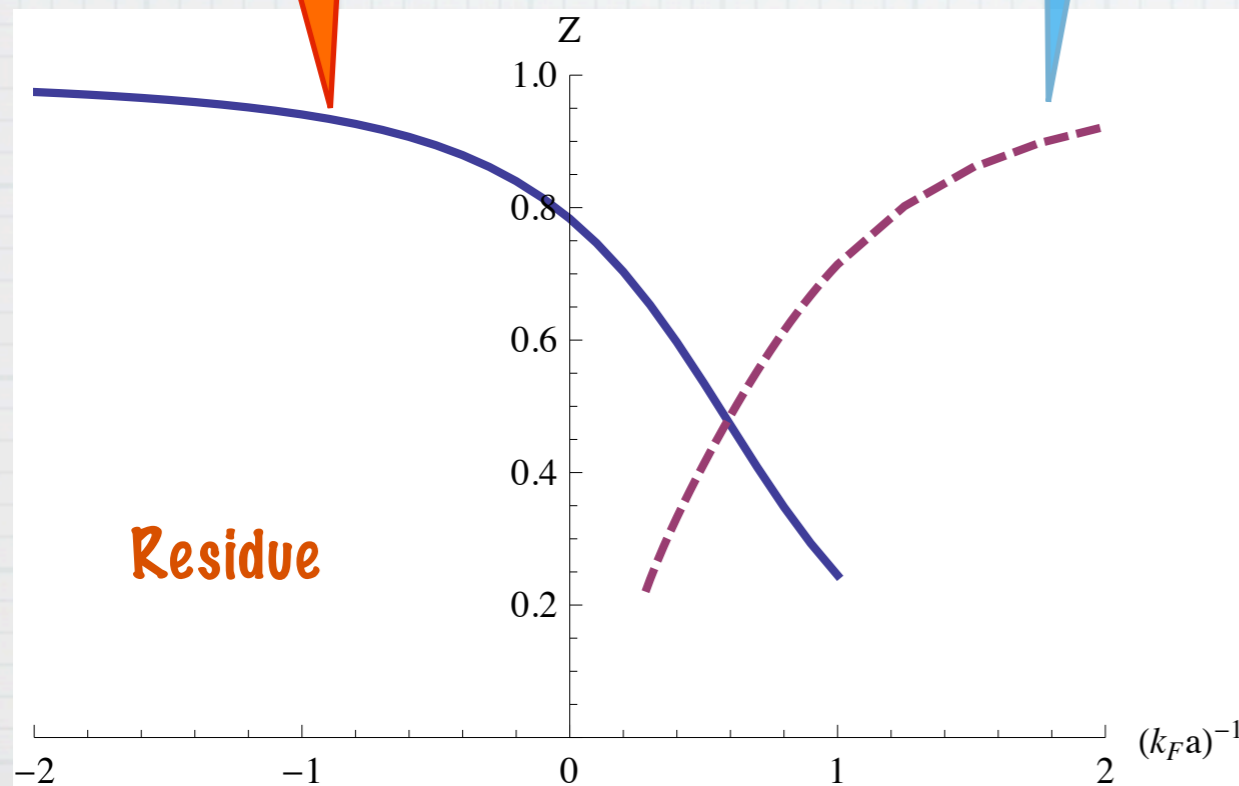
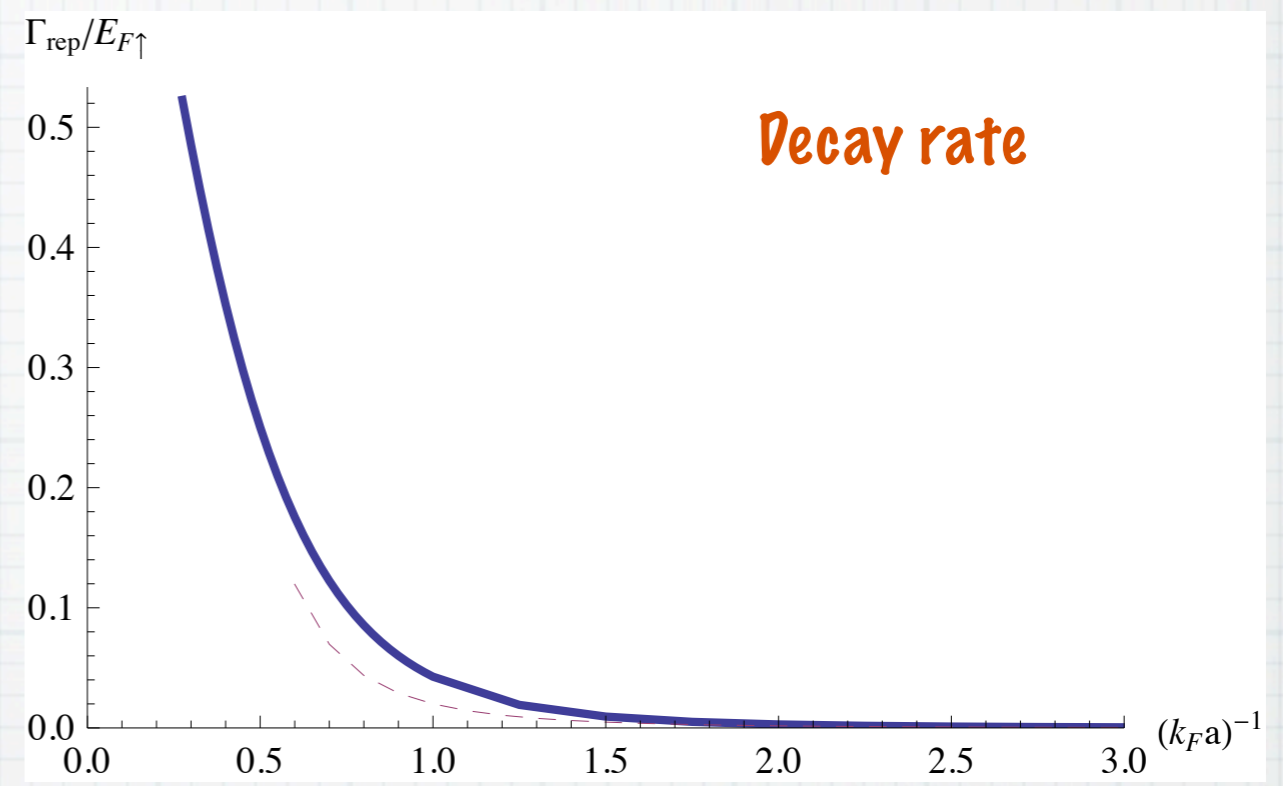
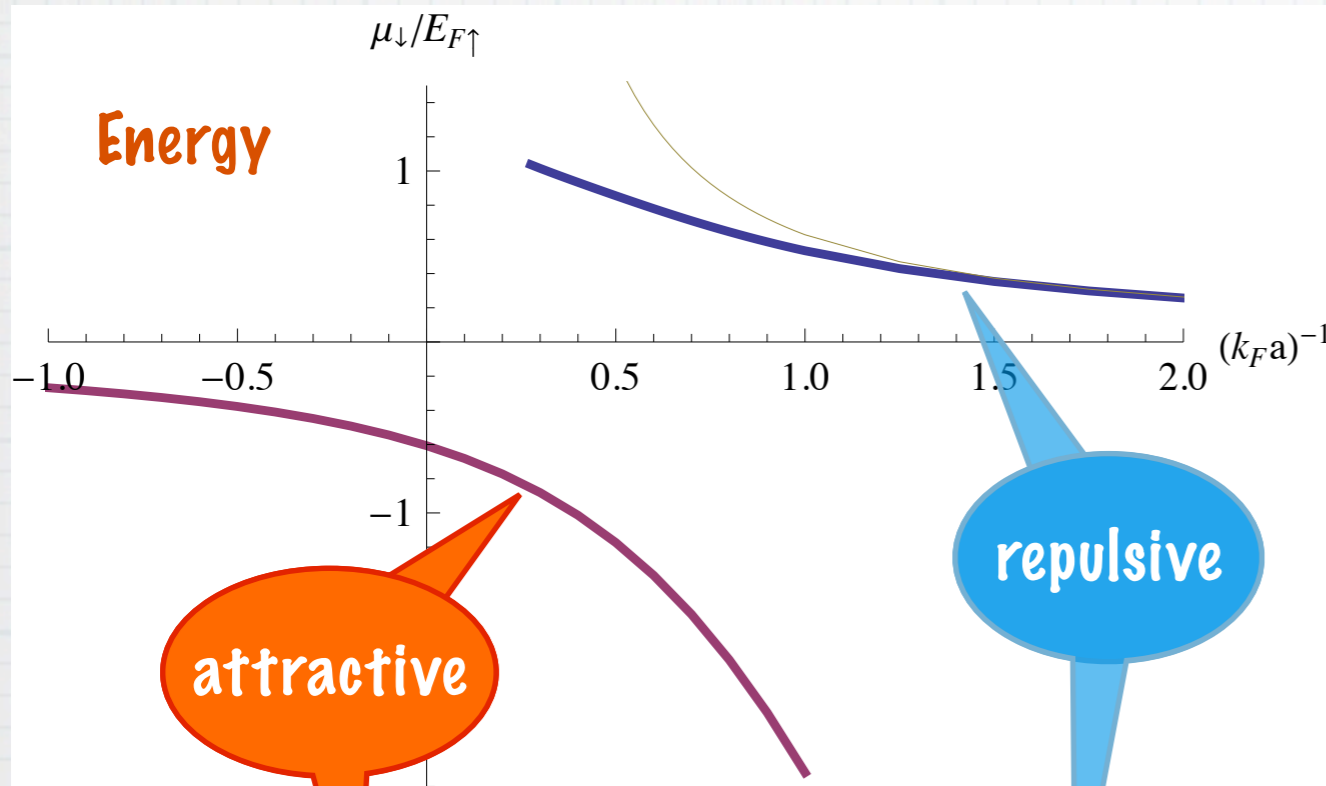
The repulsive polaron revealed



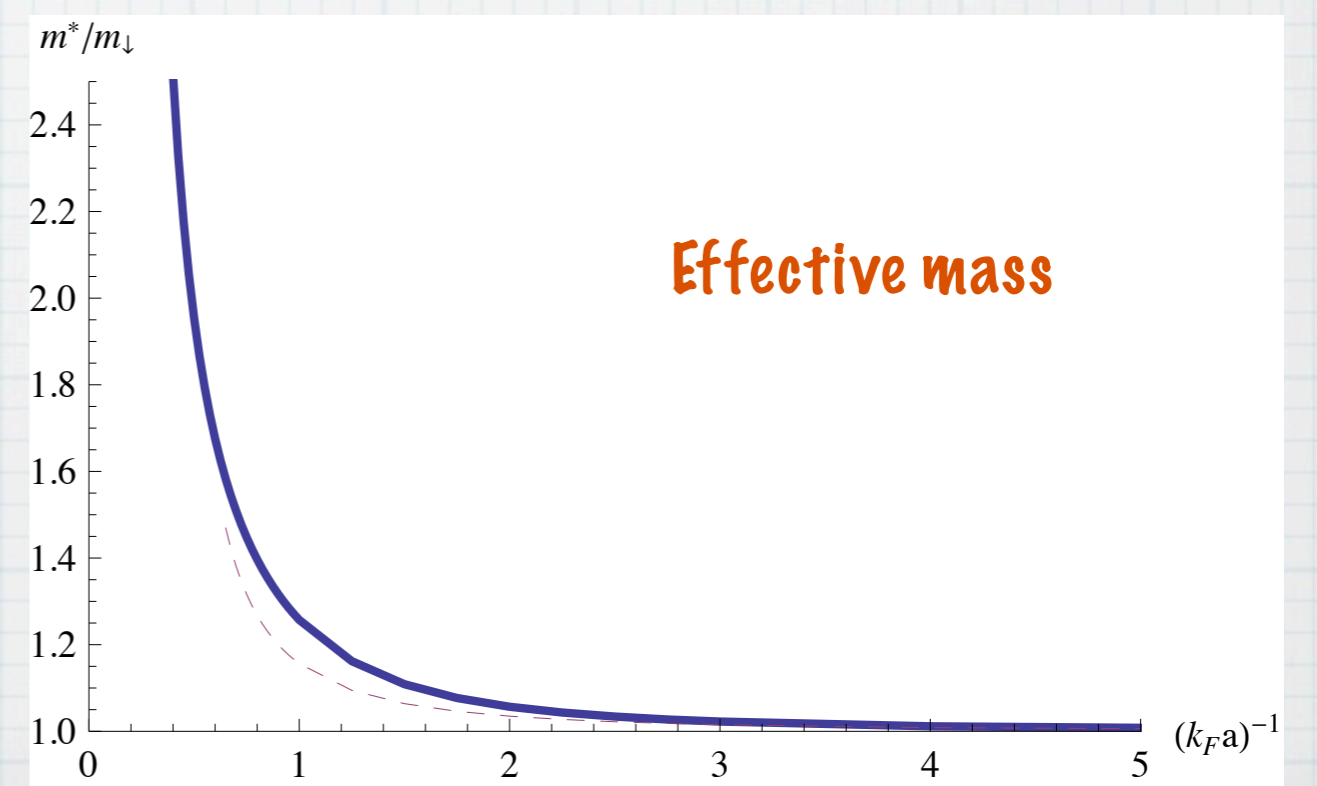
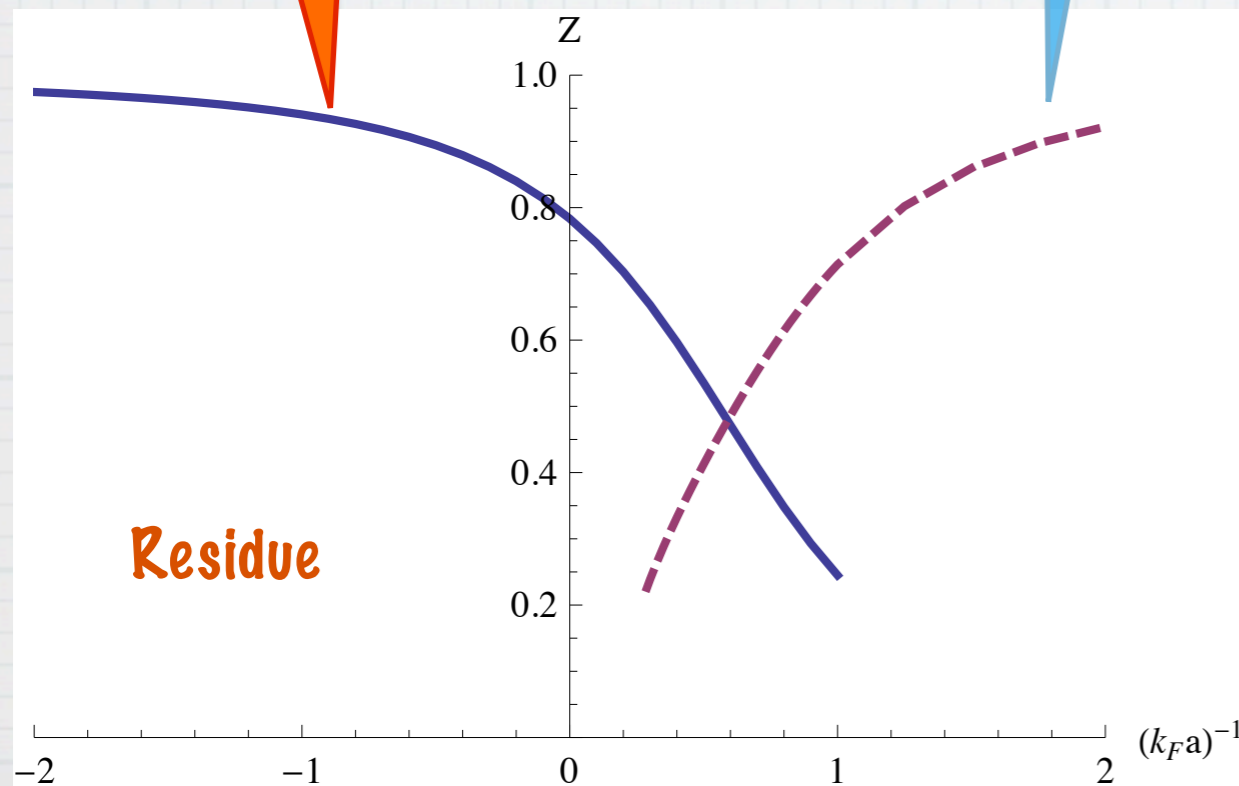
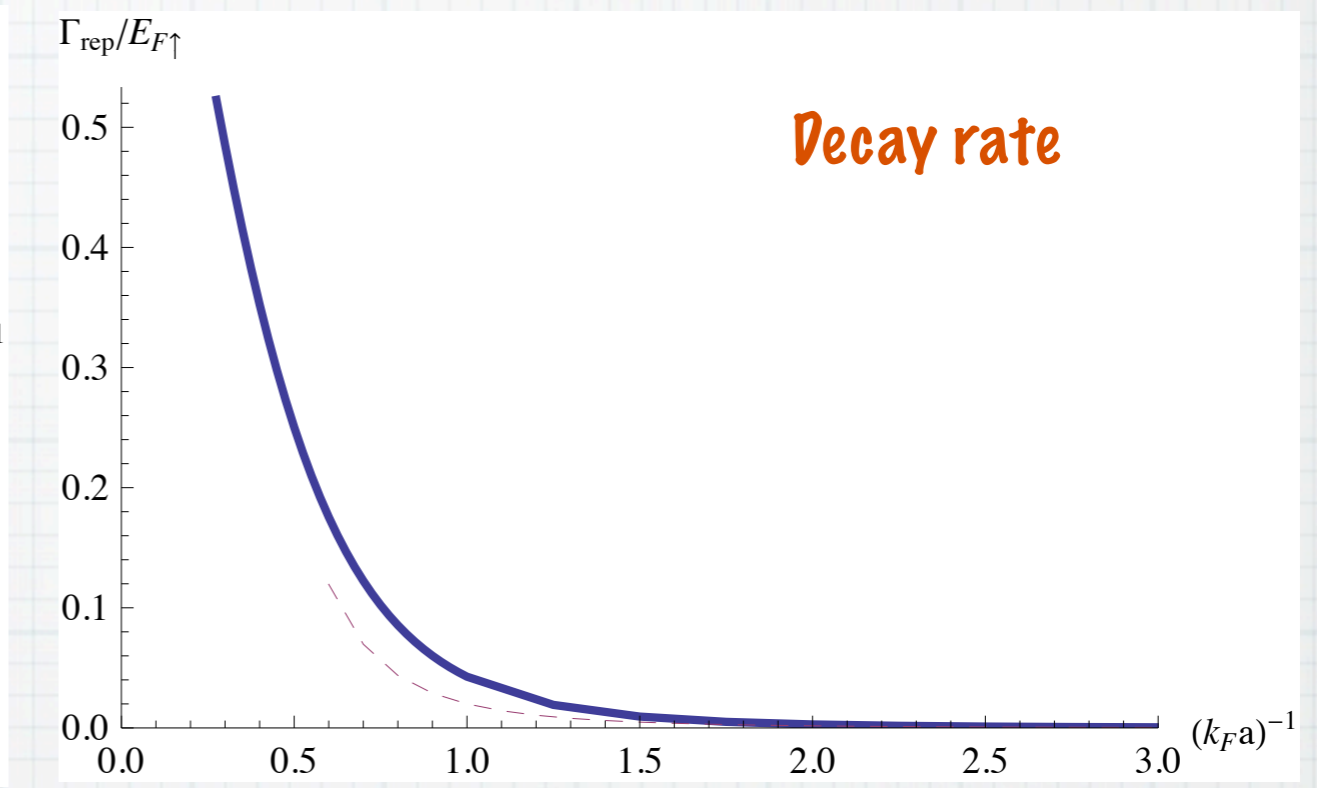
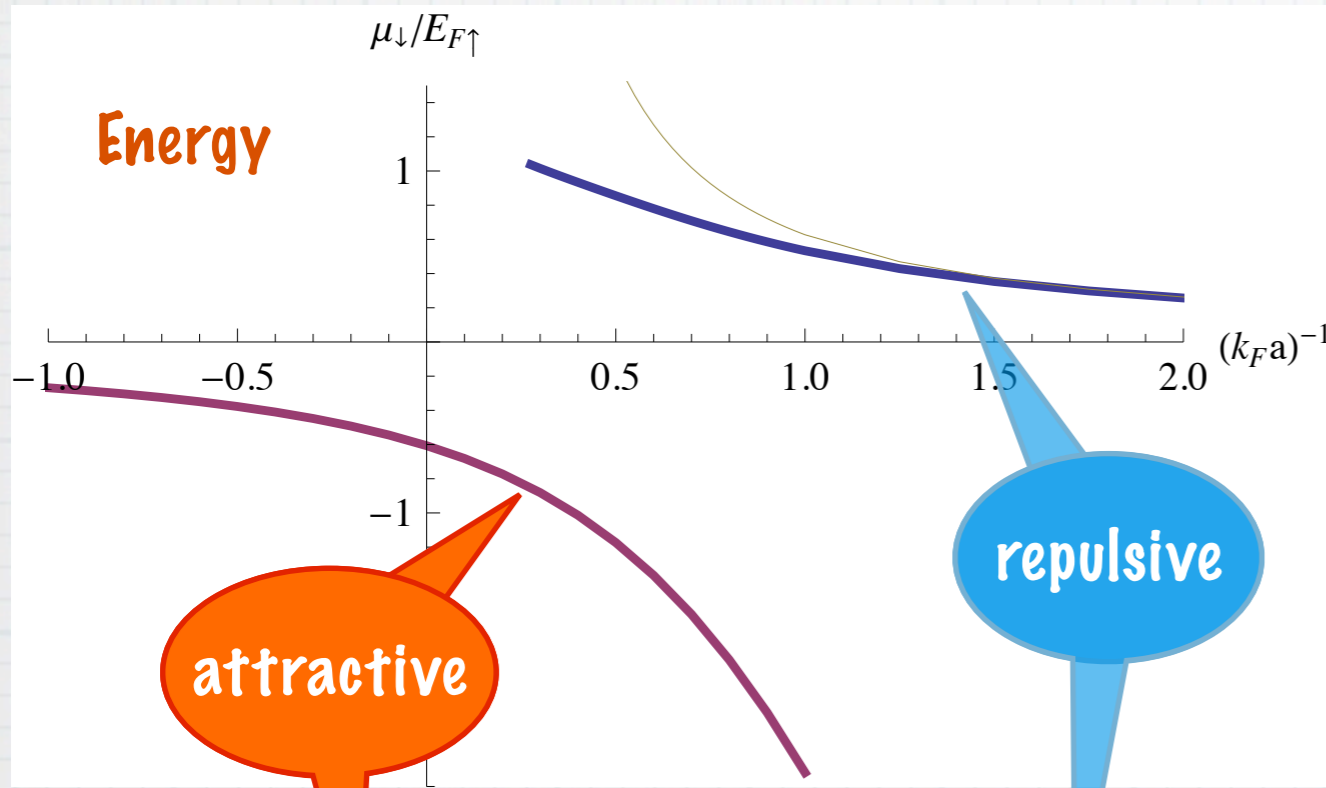
The repulsive polaron revealed



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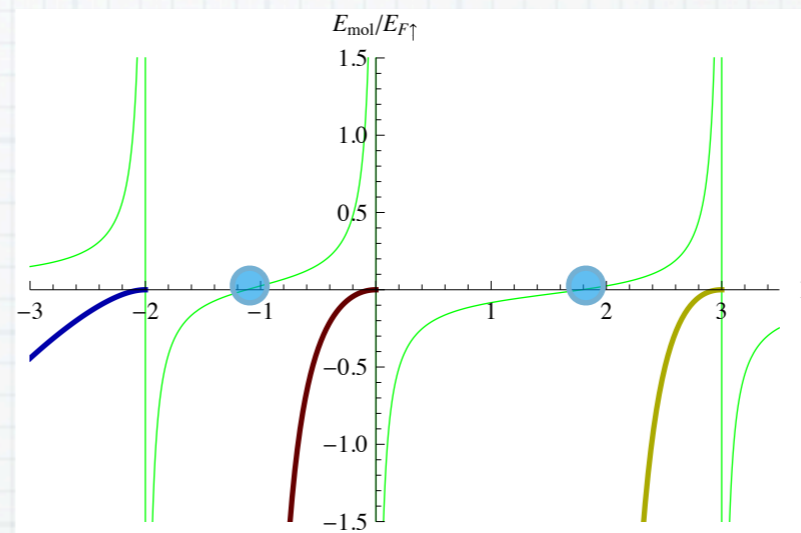


The repulsive polaron revealed



a toy model with 3 FR

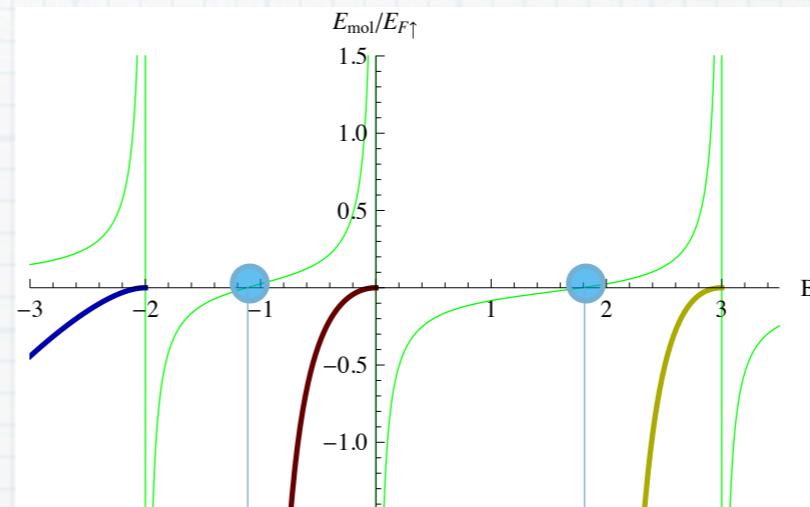
2-body bound states:



•
 $a=0$

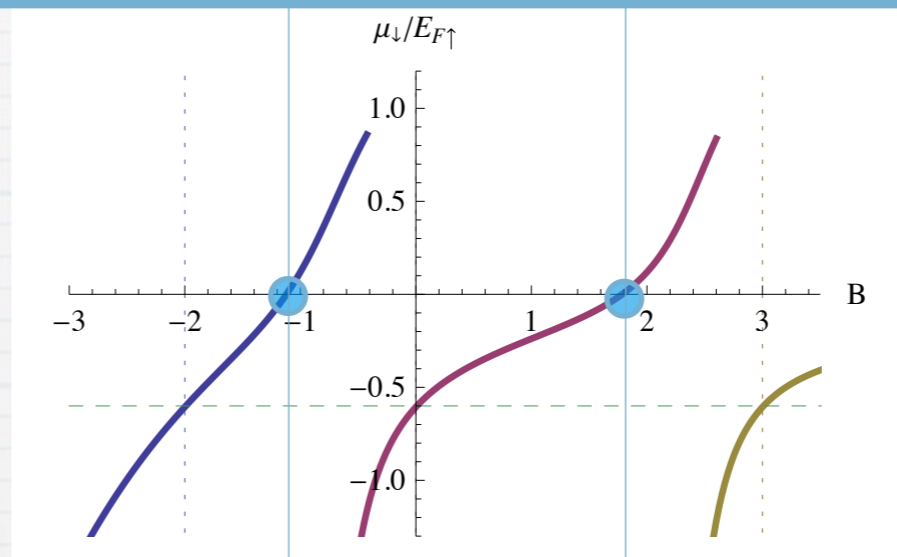
a toy model with 3 FR

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Polaronic states:

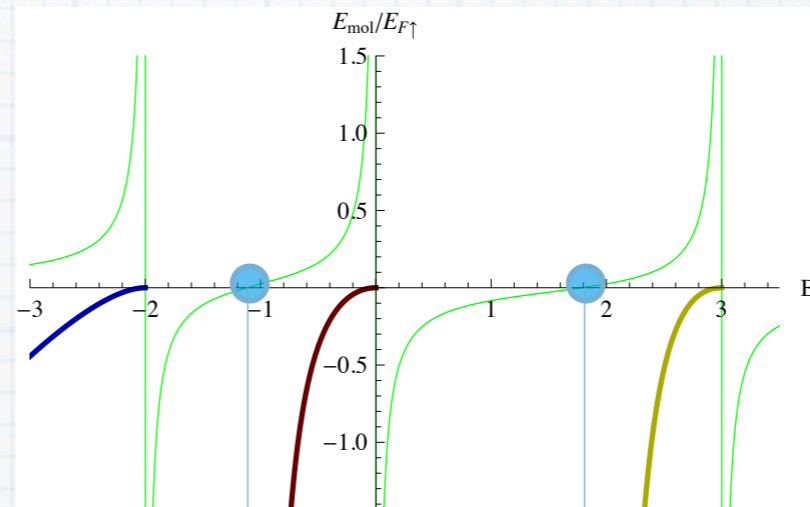


weak coupling:

$E \propto a$

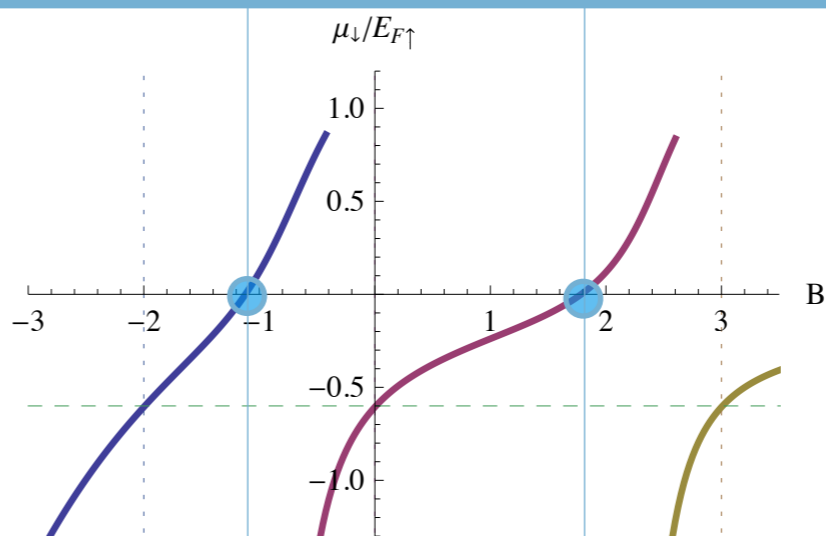
a toy model with 3 FR

2-body bound states:



•
 $a=0$

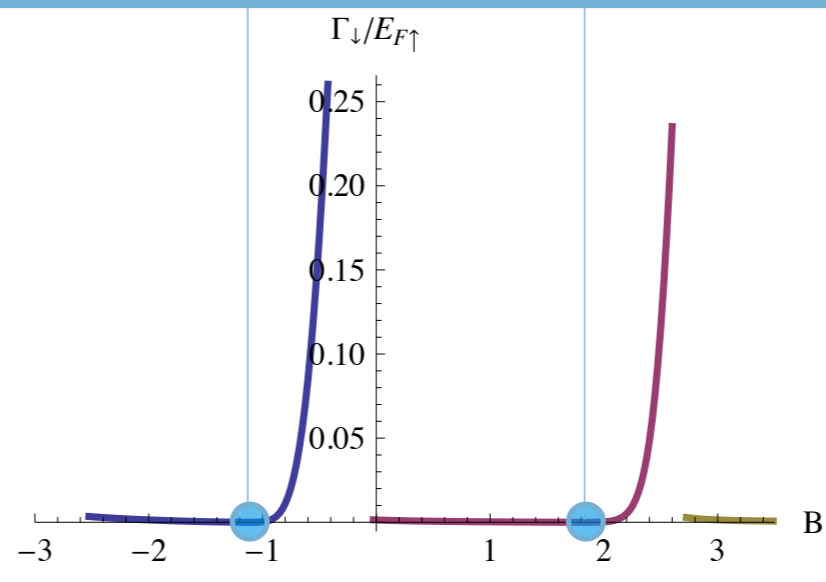
Polaronic states:



weak coupling:

$E \propto a$

Decay rates:



$\Gamma \propto \theta(a)$

Conclusions

- Quasi-particles properties fix completely the equation of state
- At small momenta, the process coupling molecules and polarons requires at least 3-bodies
- Strongly suppressed P-M decay due to a combination of small final density of states and Fermi statistics
- Expected lifetimes $\sim 10-100\text{ms}$
- Complete characterization of the repulsive branch

G. Bruun and PM, Phys. Rev. Lett. **105**, 020403 (2010)

K. Sadeghzadeh, G. Bruun, C. Lobo, PM, and A Recati, arXiv:1012.0484

PM and G. Bruun, coming soon

molecule w.f.
in vacuum:

$$\phi_q = \frac{\sqrt{8\pi a^3}}{1 + q^2 a^2} \quad \text{or} \quad \phi_r \propto \frac{e^{-r/a}}{r}$$

dressed molecule:



$$D(\mathbf{p}, \omega) \simeq \frac{Z_M}{\omega - \omega_M - p^2 / 2m_M^*}.$$

atom-molecule
coupling:

$$\bullet = \frac{1}{g(\mathbf{p}, z)} = \int \frac{d^3 q}{(2\pi)^3} \phi_q \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} \sim -\sqrt{\frac{m_r^2 a}{2\pi}}$$

(Bruun&Pethick, PRL 2004)

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(Bruun&Pethick, PRL 2004)

$$\Gamma_P = \frac{g^2 Z_M}{2} \int \frac{d^3 k d^3 q d^3 q'}{(2\pi)^9} [F(\mathbf{q}, \mathbf{k}, \omega_P) - F(\mathbf{q}', \mathbf{k}, \omega_P)]^2 \delta \left(\Delta\omega + \xi_{\mathbf{q}\uparrow} + \xi_{\mathbf{q}'\uparrow} - \xi_{\mathbf{k}\uparrow} - \frac{(\mathbf{q} + \mathbf{q}' - \mathbf{k})^2}{2m_M^*} \right)$$

$$q, q' < k_F, \quad k > k_F$$

$$F(\mathbf{q}, \mathbf{k}, \omega) = T_2(\mathbf{q}, \omega + \xi_{\mathbf{q}\uparrow}) G_{\downarrow}^0(\mathbf{q} - \mathbf{k}, \omega + \xi_{\mathbf{q}\uparrow} - \xi_{\mathbf{k}\uparrow})$$

In the neighborhood of the P-M crossing,

$$\Delta\omega \ll \epsilon_F$$
$$q \simeq k \simeq k' \simeq k_F$$

$$\int \frac{d^3k d^3q d^3q'}{(2\pi)^9} \delta(\dots) \sim (m_M^*)^{3/2} (\Delta\omega)^{7/2}$$

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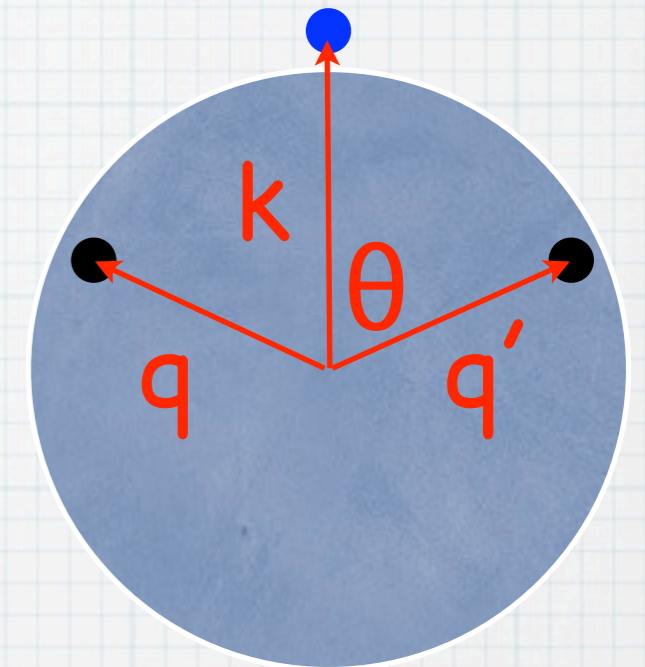
$$\Delta\omega \ll \epsilon_F$$
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The P+H+H form an equilateral triangle,
since $q + q' - k \sim 0$

At the crossing, Fermi antisymmetry
yields a vanishing of the matrix element:

$$F(\mathbf{q}, \mathbf{k}, \omega_P) - F(\mathbf{q}', \mathbf{k}, \omega_P)$$

Expand the difference to get an extra factor of $\Delta\omega$:



the angular dependence
of F is only on θ

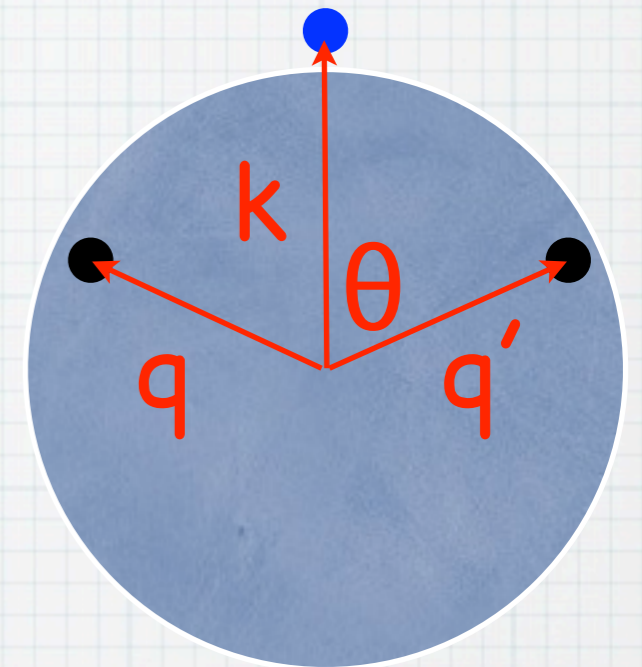
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Expand the difference to get an extra factor of $\Delta\omega$:

$$\Gamma_P \sim Z_M(k_F a) (m_M^*)^{3/2} (\Delta\omega)^{9/2}$$

1st order transition between the P&M states (no coupling at the crossing)

Mol \rightarrow Pol decay

$$\Delta\omega = \omega_P - \omega_M < 0$$

Molecule: $D(\mathbf{p}, z)^{-1} = D_0(\mathbf{p}, z)^{-1} - \Sigma_M(\mathbf{p}, z)$

Decay rate: $\Gamma_M = -\text{Im}\Sigma_M(p = 0, \omega_M)$

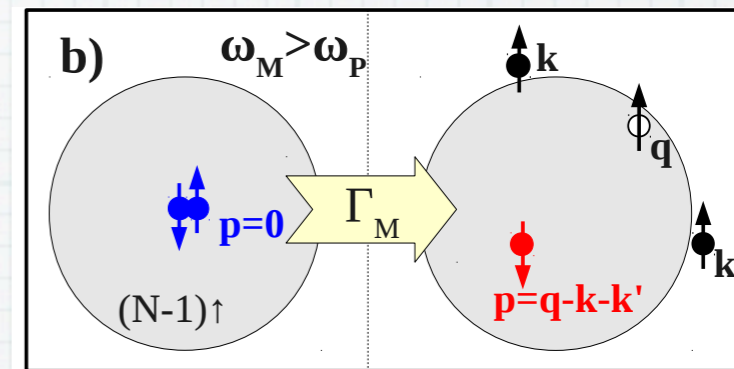
Vacuum: $D_0(\mathbf{p}, z) = \int d^3\check{q} \phi_q^2 \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} + \frac{T_2(\mathbf{p}, z)}{g(\mathbf{p}, z)^2}$

Mol \rightarrow Pol decay

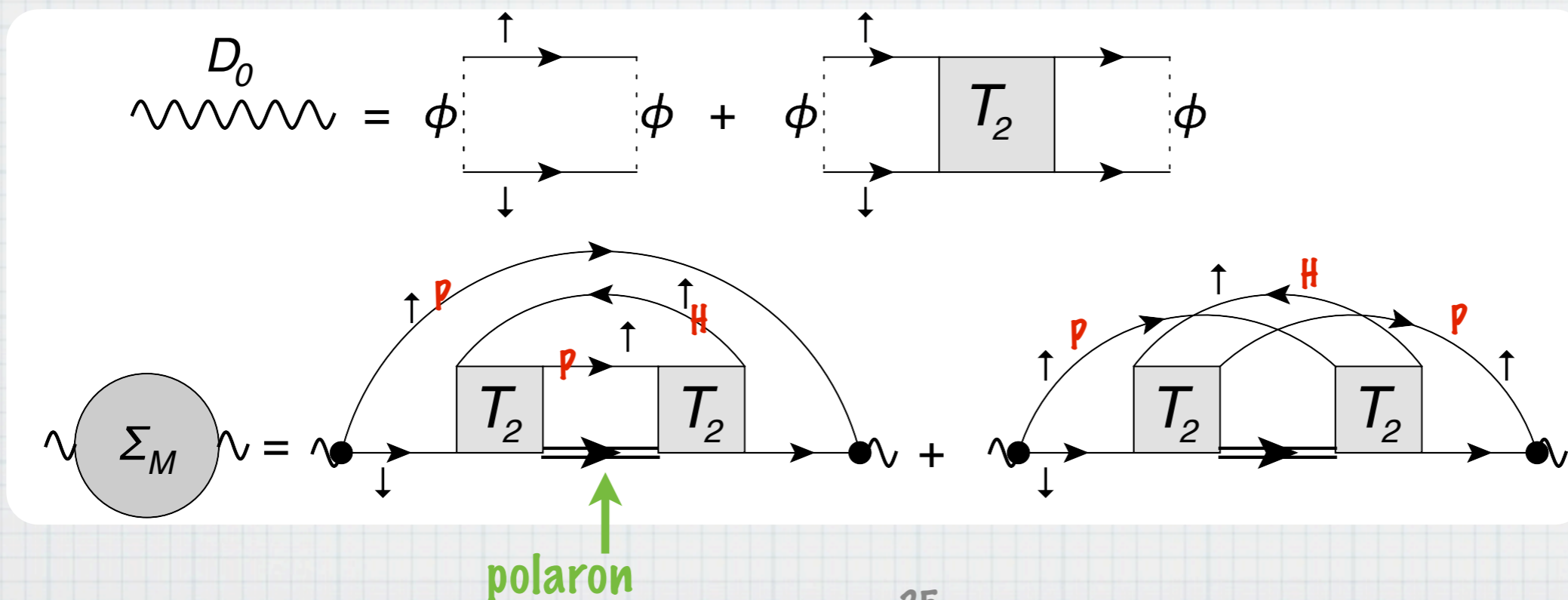
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3-body process

$$\Gamma_M = \frac{g^2 Z_P}{2} \int \frac{d^3 k d^3 k' d^3 q}{(2\pi)^9} [C(\mathbf{q}, \mathbf{k}, \omega_M) - C(\mathbf{q}, \mathbf{k}', \omega_M)]^2 \delta \left(|\Delta\omega| + \xi_{q\uparrow} - \xi_{k\uparrow} - \xi_{k'\uparrow} - \frac{(\mathbf{q} - \mathbf{k} - \mathbf{k}')^2}{2m_P^*} \right)$$

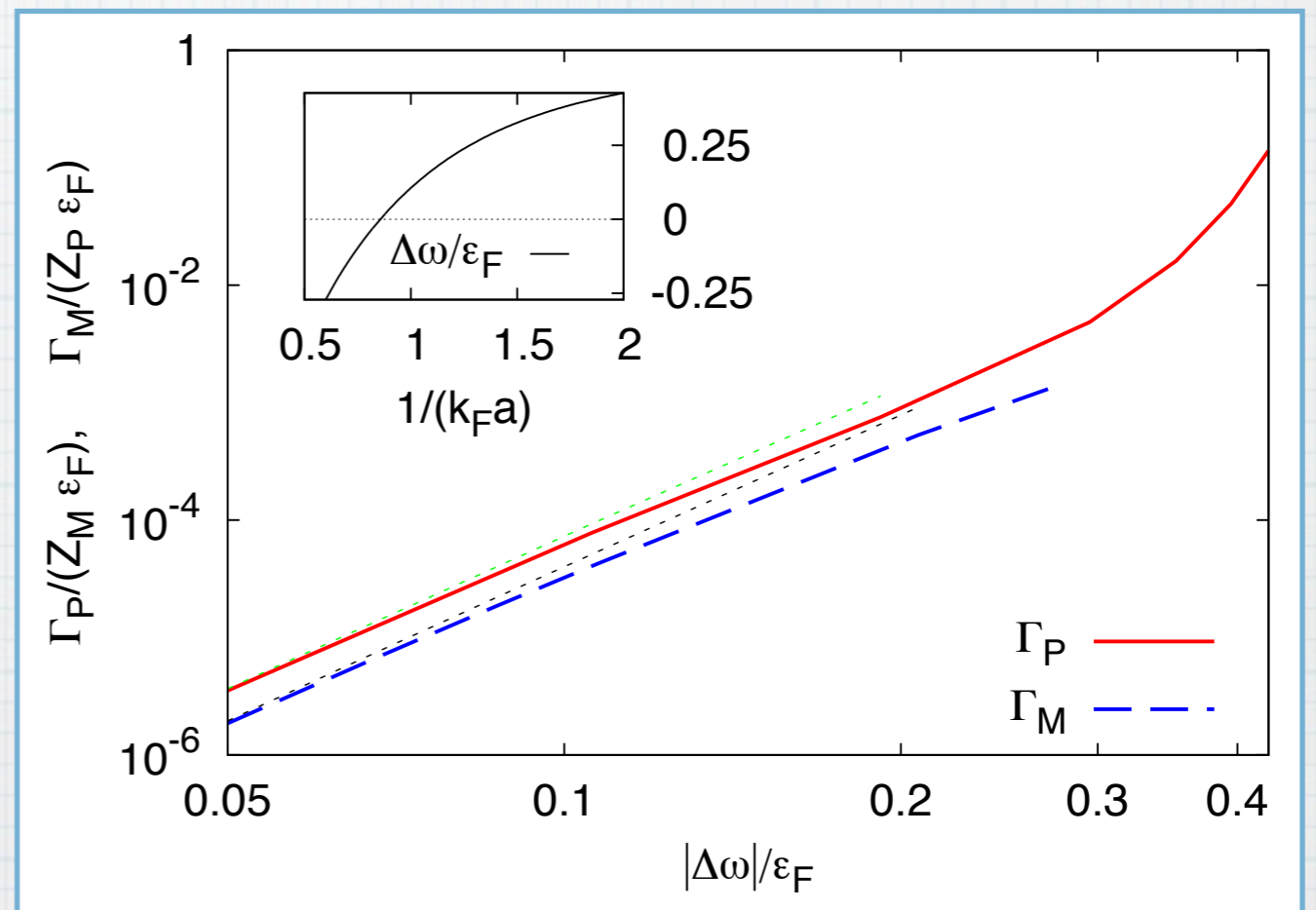
In the neighborhood of the M-P crossing, $\Gamma_M \sim Z_P(k_F a) (m_P^*)^{3/2} (-\Delta\omega)^{9/2}$

In the numerics:

$$\omega_M = -\frac{\hbar^2}{2m_r a^2} - \epsilon_F + g_3 n_\uparrow$$

$$a_3 = 1.18a$$

$$T_2(\mathbf{p}, \omega) = \frac{2\pi a/m_r}{1 - \sqrt{2m_r a^2 \left(\frac{p^2}{2m_M} - \omega - \epsilon_F + g_3 n_\uparrow \right)}}$$



$$\Gamma_M = \frac{g^2 Z_P}{2} \int \frac{d^3 k d^3 k' d^3 q}{(2\pi)^9} [C(\mathbf{q}, \mathbf{k}, \omega_M) - C(\mathbf{q}, \mathbf{k}', \omega_M)]^2 \delta \left(|\Delta\omega| + \xi_{q\uparrow} - \xi_{k\uparrow} - \xi_{k'\uparrow} - \frac{(\mathbf{q} - \mathbf{k} - \mathbf{k}')^2}{2m_P^*} \right)$$

In the neighborhood of the M-P crossing, $\Gamma_M \sim Z_P(k_F a) (m_P^*)^{3/2} (-\Delta\omega)^{9/2}$

For both decay processes,
very **long lifetimes** are ensured by:

- limited phase-space
- Fermi antisymmetry

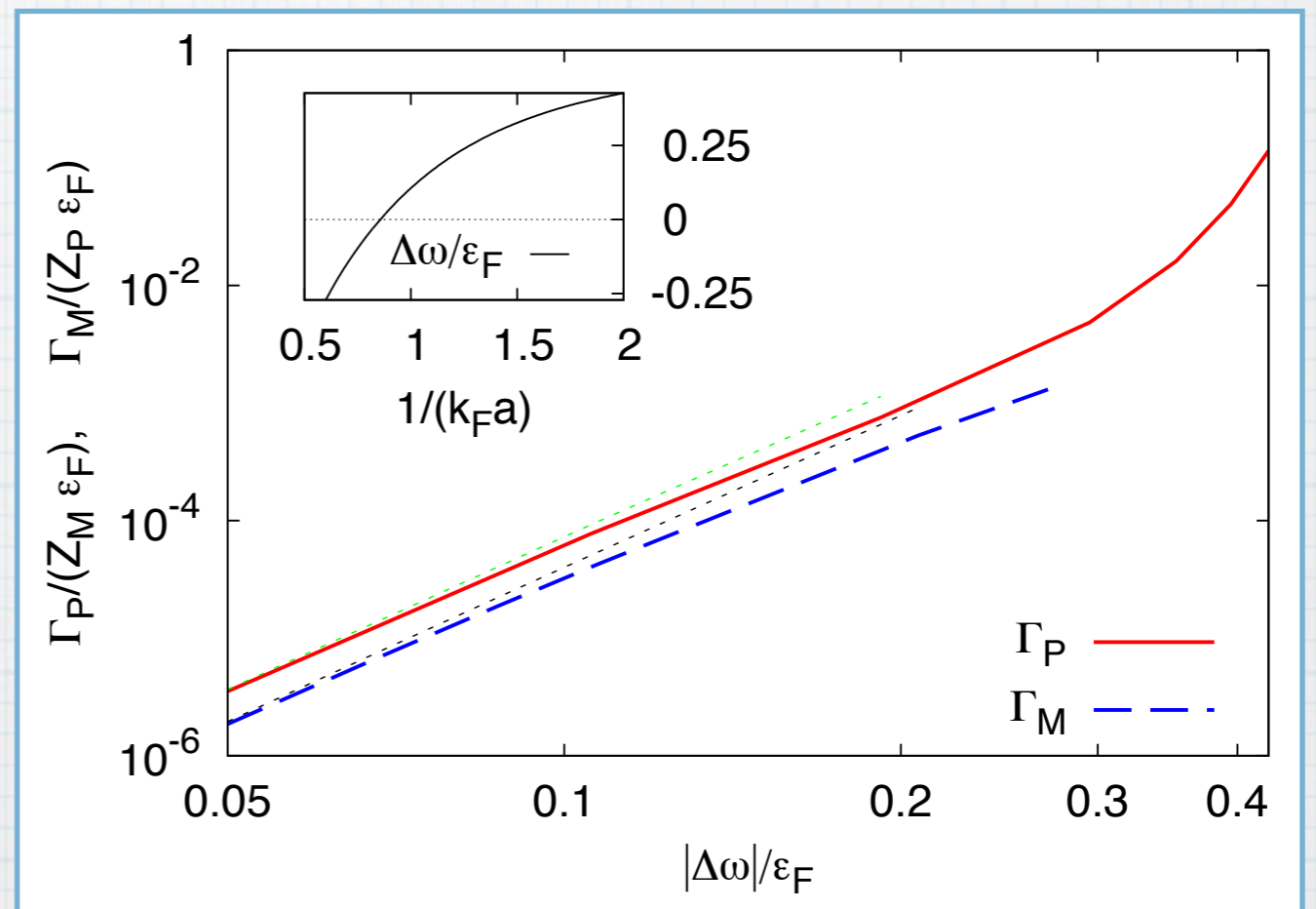
much longer than usual Fermi liquids

In the numerics:

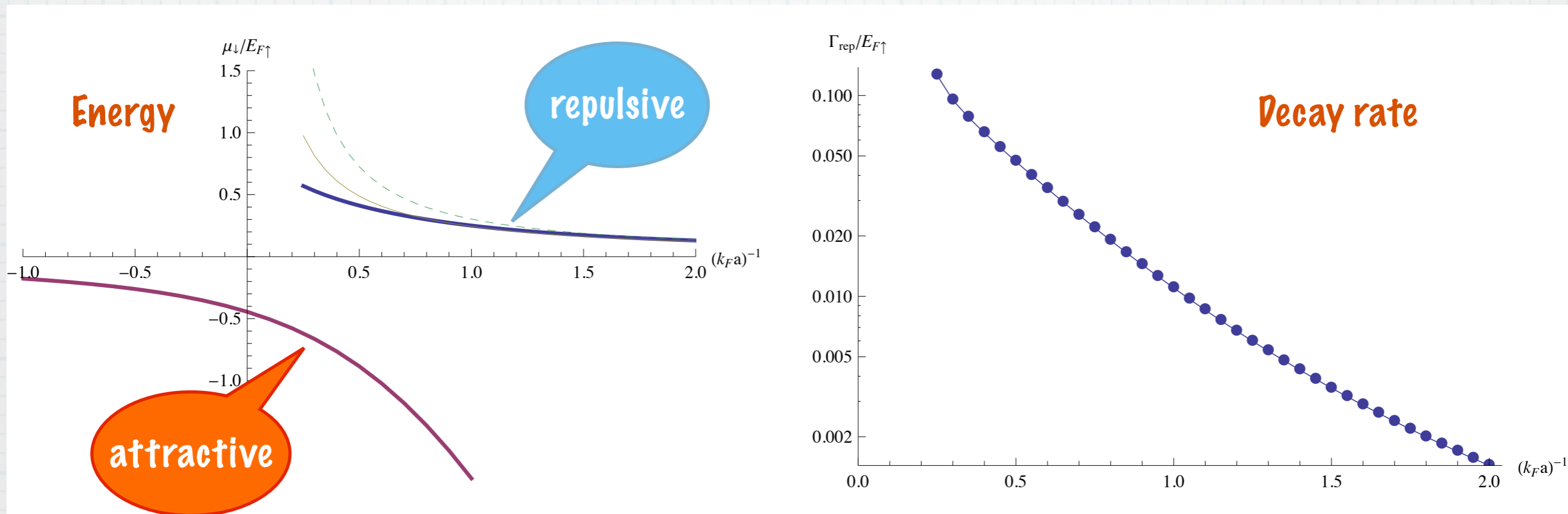
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Repulsive polaron



A ^{40}K impurity in a Fermi sea of ^6Li

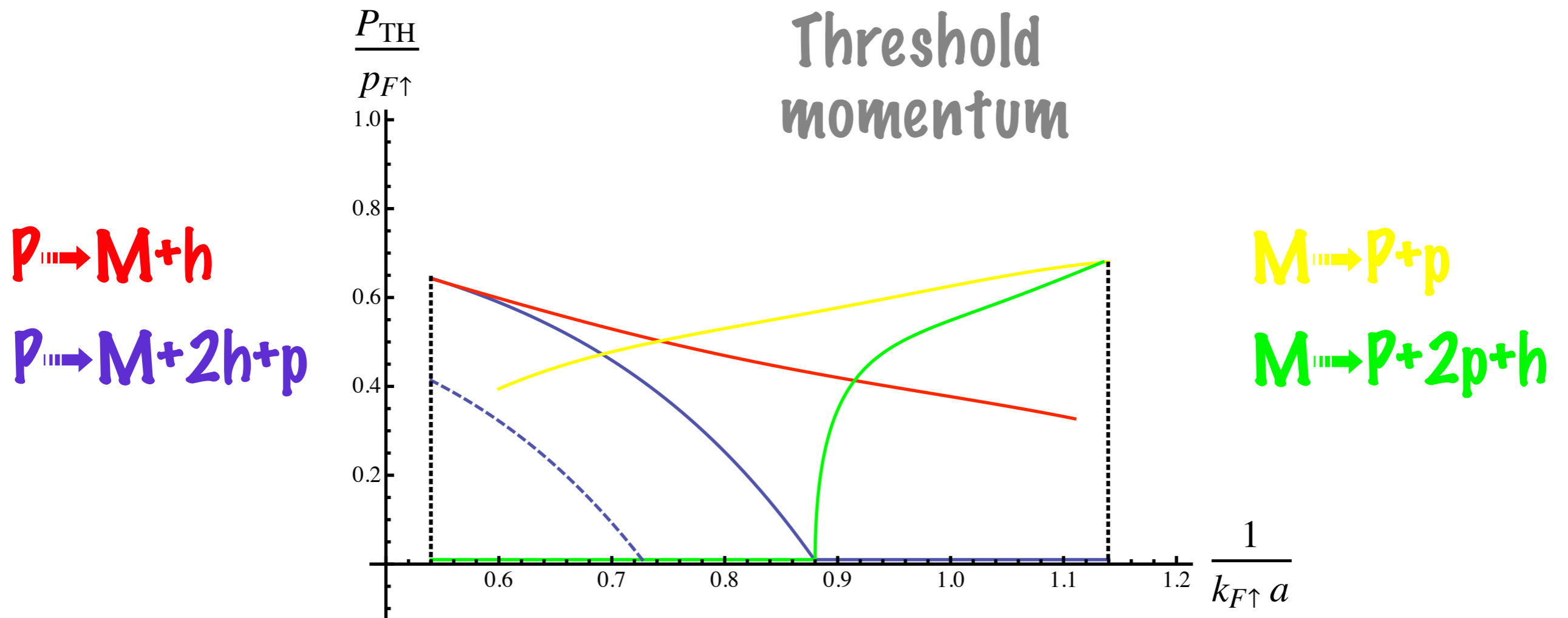
atom-molecule
coupling:

$$\bullet = \frac{1}{g(\mathbf{p}, z)} = \int \frac{d^3 q}{(2\pi)^3} \phi_q \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} \sim -\sqrt{\frac{m_r^2 a}{2\pi}}$$

Vacuum:

$$D_0(\mathbf{p}, z) = \int d^3 \check{q} \phi_q^2 \frac{1 - f(\xi_{\mathbf{p}-\mathbf{q}\uparrow})}{z - \xi_{\mathbf{p}-\mathbf{q}\uparrow} - \xi_{\mathbf{q}\downarrow}} + \frac{T_2(\mathbf{p}, z)}{g(\mathbf{p}, z)^2}$$

Decay of $p \neq 0$ QP



(preliminary)

Notice

The many-body physics discussed here can be

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defined**

Notice

The many-body physics discussed here can be

defined

calculated

Notice

The many-body physics discussed here can be

defined

calculated

and measured!